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A LIFTING-SURFACE PROGRAM FOR TRAPEZOIDAL CONTROL SURFACES WITH FLAPS

Justin E. Kerwin, et al

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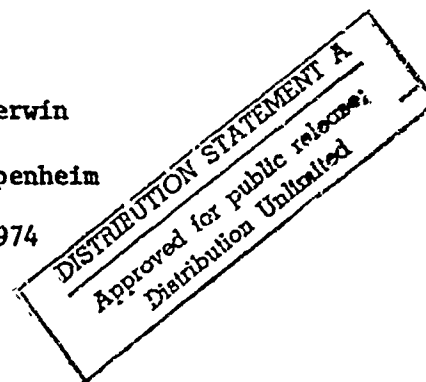
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TRAPEZOIDAL CONTROL SURFACES WITH FLAPS

by

Justin E. Kerwin  
and  
Bohdan W. Oppenheim

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# ABSTRACT

A numerical lifting surface procedure is developed specifically for flapped control surfaces with trapezoidal planforms. The procedure uses a discrete vortex approximation with spanwise vortex lines located at constant percentages of the chord. Use of the procedure in obtaining an optimum flapped rudder design is demonstrated. A listing and user's description of the computer program is included.

# NOMENCLATURE

(In the order of appearance in the text)

a	effective aspect ratio $\frac{(2s)^2}{A}$
s	semi-span
A	rudder area
f	flap area ratio $A_F/A$
$A_F$	flap area
$\lambda$	taper ratio $c_T/c_R$
$c_T$	tip chord
$c_R$	root chord
$\Lambda$	sweep angle of 1/4 chord
x	chordwise coordinate axis
z	spanwise coordinate axis
$x_L(z)$	leading edge chordwise coordinate
$x_T(z)$	trailing edge chordwise coordinate
i	spanwise panel index of the lattice
j	chordwise panel index of the lattice
m	spanwise index of control points
n	chordwise index of control points
$x_{FR}$	flap chord at the rudder root
$x_{FT}$	flap chord at the rudder tip
$x_{SR}$	skeg chord at the rudder root
$x_{ST}$	skeg chord at the rudder tip
$I_V$	spanwise precision number
$I_H$	chordwise precision number
I	number of chordwise panels on rudder
J	number of spanwise panels on rudder
NF	number of spanwise panels on flap
NS	number of spanwise panels on skeg
$G(x,z)$	nondimensional circulation distribution
$\alpha$	rudder angle of attack
$\delta$	flap deflection relative to skeg

# NOMENCLATURE (cont.)

$\gamma(x, z)$	dimensional circulation distribution
$U$	velocity at infinity
$c_{kl}$	mode amplitudes
$k$	spanwise mode index
$l$	chordwise mode index
$K$	number of spanwise modes
$L$	number of chordwise modes
$f_k(z)$	$k$ th spanwise mode
$p_l(s)$	$l$ th chordwise mode
$\tilde{z}$	transformed spanwise coordinate
$\tilde{s}$	transformed chordwise coordinate on rudder
$c(z)$	rudder local chord
$\tilde{r}$	transformed chordwise coordinate on flap
$P_l$	integral of $l$ th chordwise mode
$\Gamma(z)$	spanwise circulation distribution
$\xi, \zeta$	coordinates of a general point on a vortex
$v(s)$	velocity induced by spanwise vortex
$v(t)$	velocity induced by trailing vortex
$v_{m,n,i,j}$	velocity induced by the $(i,j)$ th element at the $(m,n)$ th control points
$v_{m,n,k,l}$	velocity induced at the $(m,n)$ th control point by the $(k,l)$ th mode of unit amplitude
$v_{m,n}$	velocity at the $(m,n)$ th control point
$\rho$	mass density
$C_{L\alpha}(z)$	local lift coefficient per unit angle of attack
$C_{L\delta}(z)$	local lift coefficient per unit flap deflection angle
$C_{L\alpha}$	overall lift coefficient per unit angle of attack
$C_{L\delta}$	overall lift coefficient per unit flap deflection angle
$C_{Di}$	induced drag coefficient
$\eta$	lifting surface efficiency
$M_l$	first moment of the $l$ th chordwise mode function
$\frac{x_H(z)}{c(z)}$	local distance of the center of pressure as a fraction of the local chord from the hinge line

NOMENCLATURE  
(cont.)

$\frac{x_{LE}(z)}{c(z)}$	local distance of the center of pressure as a fraction of the local chord from the leading edge
$\frac{x_H}{c}$	resultant chordwise position of the center of pressure relative to the flap hinge line as a fraction of the mean chord
$\bar{c}$	mean chord = $(x_{FR} + x_{FT} + x_{SR} + x_{ST})/2$
$C_{DV}$	viscous drag coefficient
$z_H$	transformed chordwise coordinate evaluated at the position of the flap hinge



## 1. INTRODUCTION

This study arose from the need to develop a rational basis for the selection of optimum geometric characteristics for rudders with relatively small flaps. In this case, very minor changes in sweep or taper can effect a major change in the spanwise distribution of flap chord. This, in turn, could be expected to influence the spanwise and chordwise distribution of lift when the flap is deflected. To optimize rudder performance one would like to have a distribution of lift which results in the maximum lift/drag ratio, highest possible stall angle, and minimum control moment for both skeg and flap. Consequently, it is necessary to have the means for estimating spanwise and chordwise distribution of lift for a given geometry.

Ship rudder effective aspect ratios typically fall in the region where neither high aspect ratio nor low aspect ratio theories are valid. One must, therefore, resort to numerical lifting-surface theory to obtain meaningful results.

It seemed most expedient for this application to write a specialized computer program designed to accommodate only trapezoidal planforms of the form shown in Fig. 1. The flap hinge is required to be at right angles to the root section, and the tip chord is required to be parallel to the flow. No restriction is placed on aspect ratio, sweep or taper, provided that the flap hinge emerges from the tip, rather than the leading or trailing edge.

A trapezoidal planform makes a discrete vortex lifting-surface model relatively simple. Unlike the original work of Faulkner [1], or current schemes for propellers [2], the present work employs spanwise vortex lines located at constant percentages of the flap and skeg chord, rather than at right angles to the oncoming flow. This eliminates the problem of vortex

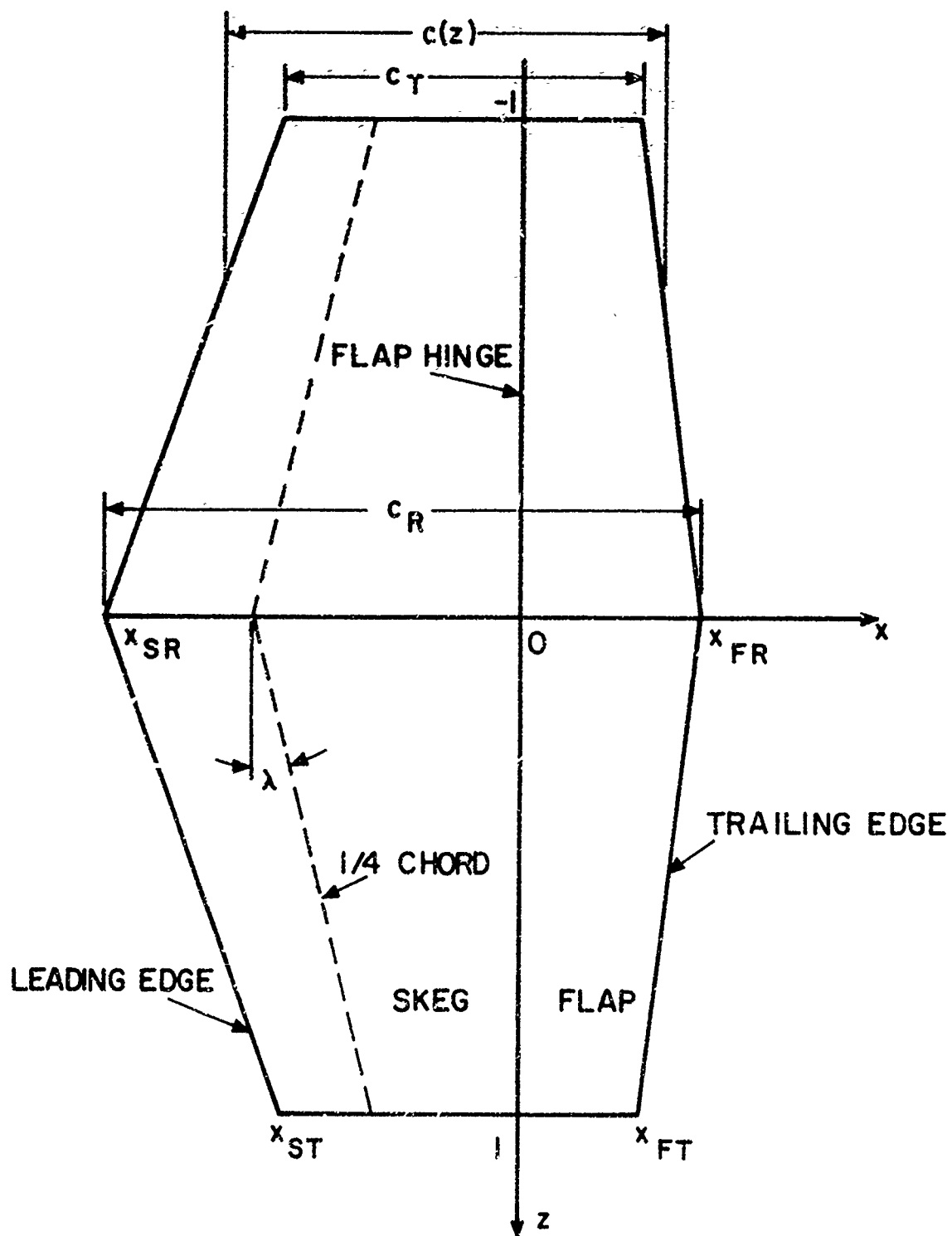


Fig. 1 Coordinate System and Notation for Planform Geometry

elements running out of the leading and trailing edges, which is inevitably a source of inaccuracy. This would be particularly objectionable in the tip region of a small flap.

As a result, the spanwise vortex lines are not purely bound vortices, since they include a component of vorticity parallel to the oncoming flow. However, as long as the accompanying system of trailing vortices is arranged in such a way that continuity of vorticity is preserved, this is an equally valid, discrete representation of a continuous vortex system.

It was decided to use a strictly linear lifting-surface theory, since an examination of results for two-dimensional airfoils with flaps [3] indicates that flap effectiveness is reduced due to viscous effects before any non-linear augment of lift becomes apparent. Consequently, the vortex system is located on a plane, even when the flap is deflected, and the influence of thickness on loading vanishes. One may then solve the problem of a rudder with zero flap deflection and unit angle of attack, and a separate problem of a rudder with zero angle of attack and unit flap deflection. The solution to any combination of angle of attack and flap deflection is then simply a linear combination of the preceding two results.

## 2. DISCRETE VORTEX ARRANGEMENT

The computer program is designed to accommodate any quadrilateral flapped control surface with a hinge axis normal to the flow, and with a tip parallel to the flow, as indicated in Fig. 1. The planform geometry is uniquely specified by the four nondimensional quantities tabulated below:

<u>Symbol</u>	<u>Definition</u>
$a$	Effective aspect ratio $\frac{(2s)^2}{A}$
$f$	Flap area ratio $A_F/A$
$\lambda$	Taper ratio $c_T/c_R$
$\Lambda$	Sweep angle of 1/4-chord

From these, we may obtain the coordinates of the four corners of the control surface. The coordinate system, as shown in Fig. 1, is located with the x-axis situated at the root section and the z-axis coincident with the flap hinge axis. The semispan,  $s$ , is taken to be unity, so that all length dimensions are nondimensionalized at the outset with respect to this quantity. The x-coordinates of the four corner points then become:

$$\begin{aligned}
 x_{FR} &= \frac{3(1 - \lambda) + 4f(1 + \lambda)}{2a(\lambda + 1)} - \frac{1}{2} \tan \Lambda \\
 x_{SR} &= x_{FR} - \frac{4}{a(\lambda + 1)} \\
 x_{FT} &= \frac{3(\lambda - 1) + 4f(1 + \lambda)}{2a(\lambda + 1)} + \frac{1}{2} \tan \Lambda \\
 x_{ST} &= x_{FT} - \frac{4\lambda}{a(\lambda + 1)} \quad .
 \end{aligned} \tag{2.1}$$

If we do not wish to have the flap hinge emerge from either the leading or trailing edge, it is necessary that the choice of input quantities be such that  $x_{FR}$  and  $x_{FT} > 0$  and  $x_{SR}$  and  $x_{ST} < 0$ . If these conditions are not met, an error message is printed.

The two trapezoidal regions representing the skeg and flap may now be subdivided into a lattice of spanwise and trailing discrete vortex lines. An individual segment of a spanwise vortex, together with the two trailing vortices originating at the ends of the segment, forms a horseshoe vortex of constant strength, as illustrated in Fig. 2.

The complete lattice arrangement is shown in Fig. 3. The fineness of the grid is controlled by specifying a spanwise precision number,  $I_V = 0, 1, \text{ or } 2$ , and a chordwise precision number,  $I_H = 0, 1, \text{ or } 2$ . A zero precision number denotes the coarsest possible grid spacing, which is the one illustrated in Fig. 3.

The semispan is divided into  $5I_V + 7$  chordwise panels of equal width. The panel nearest to the tip is further subdivided into 14 equal intervals, and this pattern is duplicated on the image side of the planform. The total number of chordwise strips over the span,  $I$ , for each precision number is as follows:

Precision Number	Number of Chordwise Panels
$I_V$	$I$
0	40
1	50
2	60

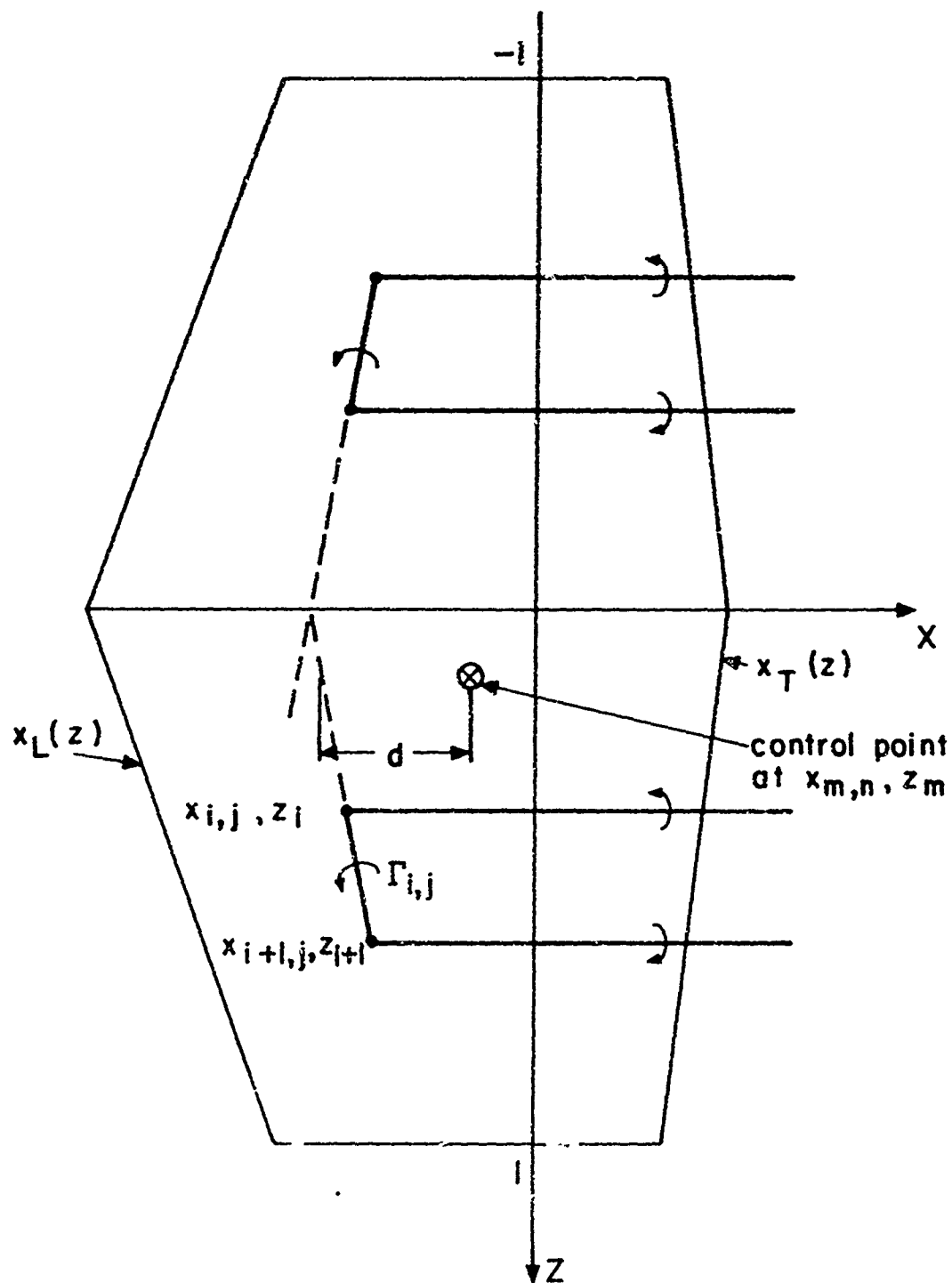


Fig. 2 Illustration of a Typical Horseshoe Vortex Element

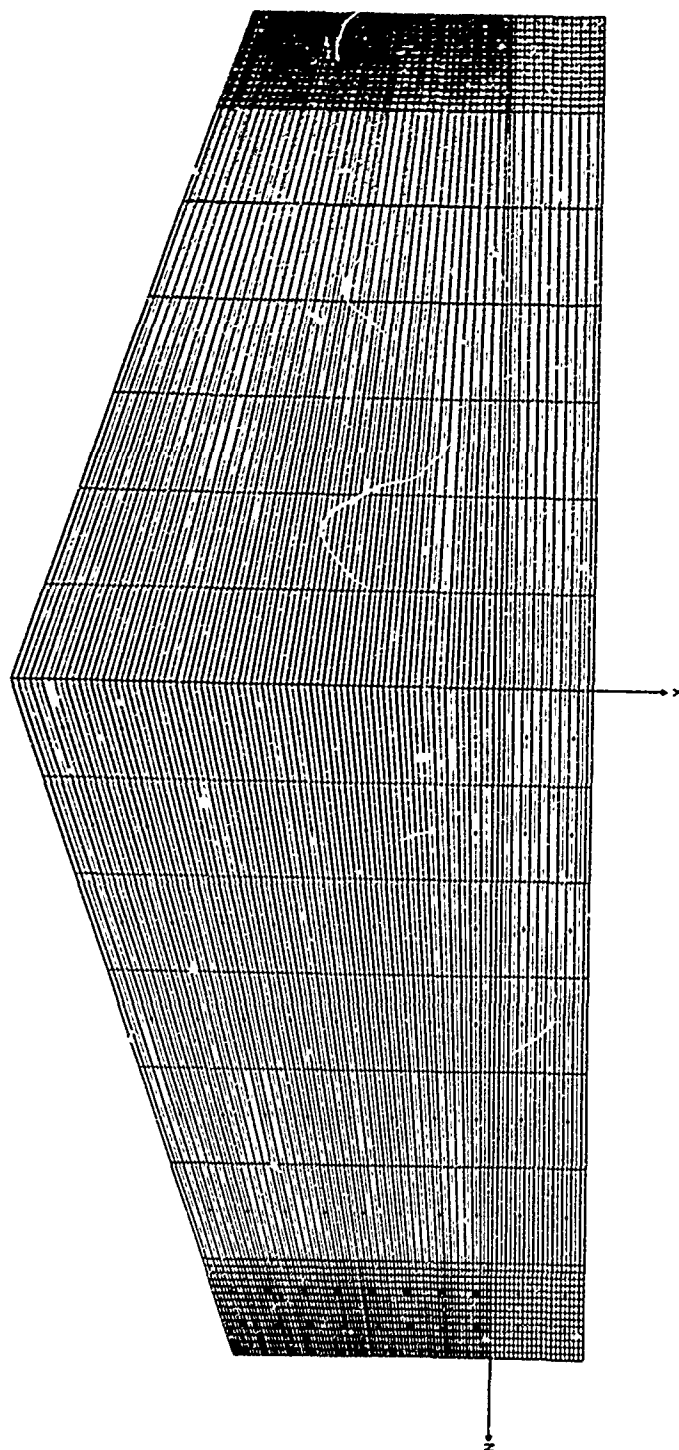


Fig. 3 Complete Lattice Arrangement Corresponding to Zero Precision Number

The total chord is divided into J spanwise panels in accordance with the specified horizontal precision number  $I_H$ ,

$$J = 50 + 10 I_H . \quad (2.2)$$

In general, it is not possible to have the chordwise spacing the same on both the skeg and the flap. However, the two spacings will be very nearly equal if the number of panels on the flap, NF, is the integer closest to the flap area ratio times the total number of intervals over the chord,

$$NF \geq f \cdot J . \quad (2.3)$$

The number of panels on the skeg, NS, is then:

$$NS = J - NF . \quad (2.4)$$

The determination of NF and NS is internal to the computer program, based on the specified input value of f.

Velocities are computed at 80 control points distributed over the span and chord. Their spanwise placement is permanent, as indicated by small circles in Fig. 3. Their chordwise placement is arbitrary, and is specified as input data by the user. Spanwise vortex lines are placed in the middle of each panel. Each control point is located midway between adjacent spanwise and trailing vortex elements, which is essential in order for the discrete system to converge to the Cauchy-Principal Value of the corresponding continuous singular integral. Thus the control points are placed on the panel boundaries.

To avoid errors due to edge effects, control points should be located at least 2 spaces away from the leading and trailing edges, and 4 spaces away from the tip. The choice of 8 spanwise and 10 chordwise stations follows from our experience with similar computing schemes for propellers [2]. Since the flap, when deflected, acts to some extent as an independent lifting



surface, it is important to include as many control points as possible over the flap chord, while avoiding the immediate vicinity of the hinge line. This requires an unusually large number of panels over the chord, and a flexible system of specifying control point locations.

This lattice arrangement, which concentrates at least half of the elements over the outer seventh of the span, is the result of considerable numerical experimentation. Lifting surfaces with trapezoidal planforms may, under certain circumstances, have a spanwise load distribution which falls to zero very abruptly at the tip. To obtain accurate results it is necessary to locate a sufficient number of control points very close to the tip. However, to avoid edge effects, as we have noted before, it is essential to have several grid spaces between the outboard control point and the tip. One must, therefore, have an extremely fine grid in this region. Since a relatively coarse spacing provides ample accuracy in the midspan region of the lifting surface, it would be an extreme waste of computer time to extend the fine grid over the entire planform.

### 3. CONTINUOUS AND DISCRETE CIRCULATION DISTRIBUTION

One could, in principle, solve for the circulation of each discrete spanwise vortex element necessary to satisfy the flow tangency boundary condition at each control point. However, this would require that the number of control points be equal to the number of vortex elements, and would result in the need to solve large systems of simultaneous equations. For zero precision number, for example, there would have to be 800 control points rather than 80, and there would be 800 simultaneous equations to solve.

We will, therefore, employ a modal approach, which greatly reduces the number of unknowns. The continuous distribution of circulation over the span and chord is assumed to be given by a series of known forms with unknown coefficients, following the classical work of Glauert [4]. The non-dimensional circulation distribution due to the angle of attack,  $G^{(\alpha)}$ , and flap deflection,  $G^{(\delta)}$ , is assumed to be given by the following series:

$$G^{(\alpha)}(x,z) = \frac{\gamma^{(\alpha)}(x,z)}{4U} = \sum_{k=1}^K \sum_{\ell=1}^{L-1} c_{k\ell}^{(\alpha)} f_k(\tilde{z}) p_{\ell}(\tilde{s}) \quad (3.1)$$

$$G^{(\delta)}(x,z) = \frac{\gamma^{(\delta)}(x,z)}{4U} = \sum_{k=1}^K \sum_{\ell=1}^L c_{k\ell}^{(\delta)} f_k(\tilde{z}) p_{\ell}(\tilde{s}) . \quad (3.2)$$

In the above equations  $c_{k\ell}$  are unknown mode amplitudes,  $f_k(\tilde{z})$  are the spanwise modes, and  $p_{\ell}(\tilde{s})$  are the chordwise modes. The spanwise modes are given by the following expression:

$$f_k(\tilde{z}) = \sin[(2k-1)\tilde{z}] , \quad (3.3)$$

where  $\tilde{z}$  is the transformed spanwise coordinate

$$\tilde{z} = \cos^{-1}(-z) . \quad (3.4)$$

The chordwise modes are

$$p_1(\tilde{s}) = \frac{2}{\pi c(z)} \left( \frac{1 + \cos \tilde{s}}{\sin \tilde{s}} \right) \quad (3.5)$$

$$p_\ell(\tilde{s}) = \frac{4 \sin[(\ell-1)\tilde{s}]}{\pi c(z)} \quad \ell = 2, 3, \dots, L-1 \quad (3.6)$$

$$p_L(\tilde{t}) = \frac{2}{\pi x_T(z)} \left( \frac{1 + \cos \tilde{t}}{\sin \tilde{t}} \right) \quad (3.7)$$

$$\text{where } \tilde{s} = \cos^{-1} \left[ 1 - \frac{2(x - x_L(z))}{c(z)} \right] \quad (3.8)$$

$$\tilde{t} = \cos^{-1} \left[ 1 - \frac{2x}{x_T(z)} \right] . \quad (3.9)$$

The last chordwise mode, corresponding to  $\ell = L$ , contains a square-root singularity at the leading edge of the flap and is, therefore, omitted in (3.1). The first six spanwise and chordwise modes are plotted in Fig. 4. One can readily see from Fig. 4 that control points must be located very near the tip in order to resolve the fifth and sixth spanwise modes.

The chordwise modes may be integrated to obtain the total circulation around any element of chord length:

$$P_1(\tilde{s}) = \int p_1(\tilde{s}) dx = \frac{1}{\pi} (\tilde{s} + \sin \tilde{s}) \quad (3.10)$$

$$P_2(\tilde{s}) = \int p_2(\tilde{s}) dx = \frac{1}{\pi} \left( \tilde{s} - \frac{\sin 2\tilde{s}}{2} \right) \quad (3.11)$$

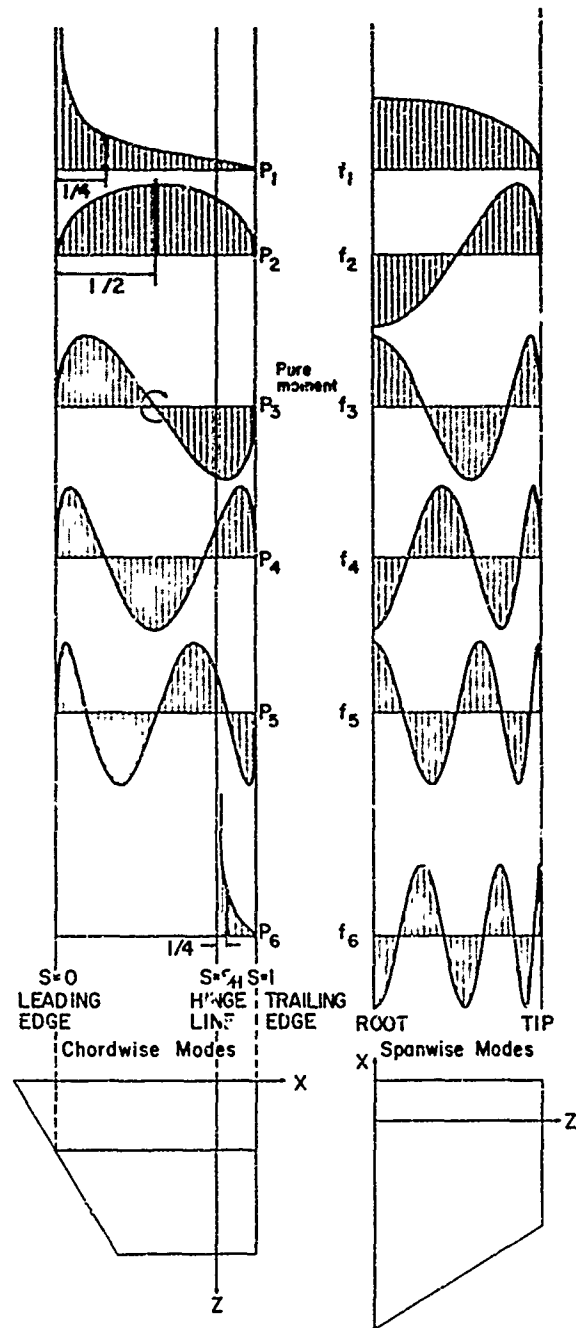


Fig. 4 Chordwise and Spanwise Mode Shapes

$$p_{\ell}(\tilde{s}) = \int p_{\ell}(\tilde{s}) dx = \frac{2}{\pi} \left( \frac{\sin(\ell-2)\tilde{s}}{2(\ell-2)} - \frac{\sin \ell \tilde{s}}{2\ell} \right) \quad \ell = 2, 3, \dots, L-1 \quad (3.12)$$

$$p_L(\tilde{t}) = \int p_L(\tilde{t}) dx = \frac{1}{\pi} (\tilde{t} + \sin \tilde{t}) \quad (3.13)$$

The total circulation around any spanwise station due to each mode is obtained by substituting the leading and trailing edges as limits in the preceding integrals. It, therefore, follows that

$$\int_{x_L(z)}^{x_T(z)} p_1(\tilde{s}) dx = 1$$

$$\int_{x_L(z)}^{x_T(z)} p_2(\tilde{s}) dx = 1$$

$$\int_{x_L(z)}^{x_T(z)} p_{\ell}(\tilde{s}) dx = 0 \quad \ell = 2, \dots, L-1$$

$$\int_0^{x_T(z)} p_L(\tilde{t}) dx = 1 \quad (3.14)$$

We can now combine (3.14) with the spanwise modes to obtain an expression for the spanwise circulation distribution

$$\Gamma^{(\alpha)}(z) = \int_{x_L(z)}^{x_T(z)} \gamma^{(\alpha)}(x, z) dx = 4U \sum_{k=1}^K (c_{k1}^{(\alpha)} + c_{k2}^{(\alpha)}) f_k(\tilde{z}) \quad (3.15)$$

$$\Gamma^{(\delta)}(z) = \int_{x_L(z)}^{x_T(z)} \gamma^{(\delta)}(x,z) dx = 4U \sum_{k=1}^K (c_{k1}^{(\delta)} + c_{k2}^{(\delta)} + c_{kL}^{(\delta)}) f_k(\tilde{z}) . \quad (3.16)$$

We can now obtain the circulation strengths of each discrete spanwise element corresponding to each mode of unit amplitude. The circulation must be, of course, constant over the span of the element. We chose to make this value correspond to that of the continuous distribution at the midspan of the element. Consequently, the spanwise mode value is obtained by substituting the value of  $\tilde{z}$  corresponding to the midspan of the element in question into (3.3). The circulation is then obtained by multiplying this value from (3.5) by the integral of the particular chordwise mode over one chordwise interval from (3.10) - (3.13). The sum of the strengths of the spanwise vortices over any chordwise panel will, therefore, be equal to the total circulation around the midspan of the panel in the continuous case.

The circulation of the two trailing vortices shed from the ends of each spanwise vortex segment must then have the same magnitude with an appropriate algebraic sign. Finally, since corresponding elements on the image portion of the lifting surface have the same strength, this computation needs to be made only over the semispan.

An alternative procedure for determining individual vortex strengths is a method originated by Faulkner [1]. The relative vortex strengths corresponding to any chordwise mode are obtained by requiring that the downwash at each panel boundary be exact in two-dimensional flow. The results are given in Table 1 for 8 chordwise panels, and in Table 2 for 50 chordwise panels. The only significant differences occur in the first

# FAULKNER SOLUTION

	J=1	J=2	J=3	J=4	J=5
N= 1	0.392761	0.072036	0.126011	0.146059	0.129591
N= 2	0.183289	0.123367	0.154209	0.068072	-0.069119
N= 3	0.126892	0.147079	0.110310	-0.065429	-0.159381
N= 4	0.096130	0.157547	0.039387	-0.149702	-0.076562
N= 5	0.074768	0.157547	-0.039387	-0.148702	0.076562
N= 6	0.057678	0.147079	-0.110310	-0.065429	0.159381
N= 7	0.042297	0.123367	-0.154209	0.068072	0.069119
N= 8	0.026184	0.072036	-0.126011	0.146059	-0.129591

# MODE INTEGRATION

	J=1	J=2	J=3	J=4	J=5
N= 1	0.440596	0.072147	0.122816	0.138168	0.116676
N= 2	0.168402	0.123354	0.152848	0.068580	-0.061543
N= 3	0.118774	0.147018	0.105588	-0.062279	-0.151446
N= 4	0.090533	0.157431	0.039160	-0.144470	-0.071452
N= 5	0.070323	0.157491	-0.039160	-0.144470	0.071452
N= 6	0.053698	0.147018	-0.105588	-0.062279	0.151446
N= 7	0.038158	0.123354	-0.152848	0.068580	0.061543
N= 8	0.019511	0.072147	-0.122816	0.138168	-0.116676

Table 1. Relative Vortex Strength for Eight Chordwise Intervals. The mode index is J and the vortex element index is N.

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	MODE INTEGRATION					PAULKNER'S SOLUTION				
	J = 1	J = 2	J = 3	J = 4	J = 5	J = 1	J = 2	J = 3	J = 4	J = 5
N = 1	0.179461	0.094773	0.009317	0.013415	0.016087	0.155178	0.004759	0.000932	0.013453	0.017118
N = 2	0.073479	0.019644	0.016232	0.021742	0.024799	0.070705	0.000000	0.000000	0.021427	0.024900
N = 3	0.055791	0.011079	0.016025	0.021773	0.024799	0.050606	0.015043	0.015043	0.021427	0.024900
N = 4	0.046534	0.012982	0.022817	0.025351	0.021157	0.040222	0.012982	0.022817	0.025351	0.021157
N = 5	0.040553	0.014566	0.023879	0.024597	0.016452	0.031752	0.014566	0.023879	0.024597	0.016452
N = 6	0.036257	0.015528	0.023840	0.022817	0.010760	0.021752	0.015528	0.023840	0.022817	0.010760
N = 7	0.032965	0.017122	0.025333	0.023369	0.008431	0.014764	0.017122	0.025333	0.023369	0.008431
N = 8	0.030324	0.018191	0.026447	0.021745	0.008003	0.010900	0.018191	0.026447	0.021745	0.008003
N = 9	0.028299	0.019127	0.027524	0.021454	0.007648	0.008431	0.019127	0.027524	0.021454	0.007648
N = 10	0.026203	0.019976	0.028475	0.021174	0.007216	0.007216	0.019976	0.028475	0.021174	0.007216
N = 11	0.024203	0.020741	0.029455	0.020815	0.006763	0.006763	0.020741	0.029455	0.020815	0.006763
N = 12	0.022251	0.021430	0.030466	0.020466	0.006311	0.006311	0.021430	0.030466	0.020466	0.006311
N = 13	0.020258	0.022051	0.031504	0.020147	0.005860	0.005860	0.022051	0.031504	0.020147	0.005860
N = 14	0.018260	0.022639	0.032566	0.019826	0.005410	0.005410	0.022639	0.032566	0.019826	0.005410
N = 15	0.016260	0.023198	0.033647	0.019504	0.004960	0.004960	0.023198	0.033647	0.019504	0.004960
N = 16	0.014260	0.023752	0.034731	0.019182	0.004510	0.004510	0.023752	0.034731	0.019182	0.004510
N = 17	0.012260	0.024306	0.035815	0.018860	0.004060	0.004060	0.024306	0.035815	0.018860	0.004060
N = 18	0.010260	0.024859	0.036900	0.018538	0.003610	0.003610	0.024859	0.036900	0.018538	0.003610
N = 19	0.008260	0.025413	0.037984	0.018216	0.003160	0.003160	0.025413	0.037984	0.018216	0.003160
N = 20	0.006260	0.025967	0.039068	0.017894	0.002710	0.002710	0.025967	0.039068	0.017894	0.002710
N = 21	0.004260	0.026521	0.040152	0.017572	0.002260	0.002260	0.026521	0.040152	0.017572	0.002260
N = 22	0.002260	0.027075	0.041236	0.017250	0.001810	0.001810	0.027075	0.041236	0.017250	0.001810
N = 23	0.000260	0.027629	0.042320	0.016928	0.001360	0.001360	0.027629	0.042320	0.016928	0.001360
N = 24	0.000000	0.028183	0.043404	0.016606	0.000910	0.000910	0.028183	0.043404	0.016606	0.000910
N = 25	0.000000	0.028737	0.044488	0.016284	0.000460	0.000460	0.028737	0.044488	0.016284	0.000460
N = 26	0.000000	0.029291	0.045572	0.015962	0.000010	0.000010	0.029291	0.045572	0.015962	0.000010
N = 27	0.000000	0.029845	0.046656	0.015640	0.000000	0.000000	0.029845	0.046656	0.015640	0.000000
N = 28	0.000000	0.030399	0.047740	0.015318	0.000000	0.000000	0.030399	0.047740	0.015318	0.000000
N = 29	0.000000	0.030953	0.048824	0.014996	0.000000	0.000000	0.030953	0.048824	0.014996	0.000000
N = 30	0.000000	0.031507	0.049908	0.014674	0.000000	0.000000	0.031507	0.049908	0.014674	0.000000
N = 31	0.000000	0.032061	0.050992	0.014352	0.000000	0.000000	0.032061	0.050992	0.014352	0.000000
N = 32	0.000000	0.032615	0.052076	0.014030	0.000000	0.000000	0.032615	0.052076	0.014030	0.000000
N = 33	0.000000	0.033169	0.053160	0.013708	0.000000	0.000000	0.033169	0.053160	0.013708	0.000000
N = 34	0.000000	0.033723	0.054244	0.013386	0.000000	0.000000	0.033723	0.054244	0.013386	0.000000
N = 35	0.000000	0.034277	0.055328	0.013064	0.000000	0.000000	0.034277	0.055328	0.013064	0.000000
N = 36	0.000000	0.034831	0.056412	0.012742	0.000000	0.000000	0.034831	0.056412	0.012742	0.000000
N = 37	0.000000	0.035385	0.057496	0.012420	0.000000	0.000000	0.035385	0.057496	0.012420	0.000000
N = 38	0.000000	0.035939	0.058580	0.012098	0.000000	0.000000	0.035939	0.058580	0.012098	0.000000
N = 39	0.000000	0.036493	0.059664	0.011776	0.000000	0.000000	0.036493	0.059664	0.011776	0.000000
N = 40	0.000000	0.037047	0.060748	0.011454	0.000000	0.000000	0.037047	0.060748	0.011454	0.000000
N = 41	0.000000	0.037601	0.061832	0.011132	0.000000	0.000000	0.037601	0.061832	0.011132	0.000000
N = 42	0.000000	0.038155	0.062916	0.010810	0.000000	0.000000	0.038155	0.062916	0.010810	0.000000
N = 43	0.000000	0.038709	0.064000	0.010488	0.000000	0.000000	0.038709	0.064000	0.010488	0.000000
N = 44	0.000000	0.039263	0.065084	0.010166	0.000000	0.000000	0.039263	0.065084	0.010166	0.000000
N = 45	0.000000	0.039817	0.066168	0.009844	0.000000	0.000000	0.039817	0.066168	0.009844	0.000000
N = 46	0.000000	0.040371	0.067252	0.009522	0.000000	0.000000	0.040371	0.067252	0.009522	0.000000
N = 47	0.000000	0.040925	0.068336	0.009200	0.000000	0.000000	0.040925	0.068336	0.009200	0.000000
N = 48	0.000000	0.041479	0.069420	0.008878	0.000000	0.000000	0.041479	0.069420	0.008878	0.000000
N = 49	0.000000	0.042033	0.070504	0.008556	0.000000	0.000000	0.042033	0.070504	0.008556	0.000000
N = 50	0.000000	0.042587	0.071588	0.008234	0.000000	0.000000	0.042587	0.071588	0.008234	0.000000

Table 2. Relative Vortex Strength for Fifty Chordwise Intervals. The mode index is J and the vortex element index is N.



chordwise mode near the leading and trailing edges, and both methods result in essentially identical final values of lift distribution. However, the Faulkner method requires a greater amount of computation, since the relative panel widths on the skeg and flap are functions of spanwise position. For zero precision number, the Faulkner method would require the solution of 120 different sets of simultaneous equations, each with 50 unknowns. Without the complication of the flap, the Faulkner procedure for obtaining relative vortex strengths would be much simpler, since the relative vortex strengths could be pre-computed.

#### 4. COMPUTATION OF INDUCED VELOCITIES

The lattice arrangement as described in Section 2 results in a set of discrete, skewed horseshoe vortex elements, as shown in Fig. 2. Consider a particular horseshoe element consisting of a spanwise vortex of strength  $\Gamma$  extending from  $(x_1, z_1)$  to  $(x_2, z_2)$ , joined by semi-infinite trailing vortices of strength  $\Gamma$  starting at  $(x_2, z_2)$  and  $-\Gamma$  starting at  $(x_1, z_1)$ . For each such element located in the interval  $0 \leq z \leq 1$ , there will be a corresponding image element of the same strength extending from  $(x_2, -z_2)$  to  $(x_1, -z_1)$ . The trailing vortex shed at  $(x_2, -z_2)$  will have a strength  $-\Gamma$ , while the vortex shed from  $(x_1, -z_1)$  will have a strength of  $+\Gamma$ .

Let us first consider the spanwise vortex segment. Defining  $(\xi, \zeta)$  as the coordinates of a general point on the vortex, and  $(x, z)$  as the coordinates of the control point, the velocity may be written in accordance with the law of Biot-Savart as follows:

$$\frac{4\pi v^{(s)}(x, z)}{\Gamma} = \int_{x_1, z_1}^{x_2, z_2} \frac{(x-\xi)d\zeta - (z-\zeta)d\xi}{[(x-\xi)^2 + (z-\zeta)^2]^{3/2}} \quad (4.1)$$

Along the vortex, we have:

$$t = \frac{d\xi}{d\zeta} = \frac{x_2 - x_1}{z_2 - z_1} = \text{const} \quad , \quad (4.2)$$

so that (4.1) may be expressed in terms of the variable  $\zeta$  alone and readily integrated to give the result:

$$\frac{4\pi v^{(s)}(x, z)}{\Gamma} = \frac{2a\zeta + b}{2d\sqrt{a\zeta^2 + b\zeta + c}} \Bigg|_{z_1}^{z_2}$$

where

$$\begin{aligned}
 a &= 1+t^2 \\
 e &= x-x_1+tz_1 \\
 b &= -2(et+z) \\
 c &= e^2+z^2 \\
 d &= e-tz \quad .
 \end{aligned}
 \tag{4.3}$$

Equation (4.3) is not suitable for numerical computation if  $d$  becomes small. As is evident from Fig. 2,  $d$  is the horizontal distance between the control point and the spanwise vortex line or its extension. For the vortex arrangement shown in Fig. 3,  $d$  can never be less than approximately one half the chordwise grid spacing for vortex elements occupying the semi-span interval  $0 \leq z \leq 1$ . The distance  $d$  will only become small enough to cause problems if the aspect ratio is extremely high. For this situation, the limiting form of (4.1) valid for  $d \ll 1$  is:

$$\frac{4\pi v^{(s)}(x,z)}{\Gamma} = 2\sqrt{a} \left[ \frac{1}{d} - \frac{2d}{(2az_1+b)^2} - \frac{2d}{(2az_2+b)^2} \right] , \tag{4.4}$$

provided  $z_1 < z < z_2$ . In that case, the velocity approaches that of an infinite vortex as given by the first term in (4.4). The second and third terms represent small corrections to the result for an infinite vortex. It has been found by numerical experimentation that round-off error is minimized if (4.4) is used for  $|d| < 0.002$ .

If  $z$  is not in the interval  $z_1 < z < z_2$ , the first term of (4.4) disappears, and the velocity tends to zero as  $d \rightarrow 0$ . If (4.3) is used in this case, catastrophic round-off error can occur. This is because the integral becomes very large, but independent of  $\zeta$ . Hence, the integral

from  $z_1$  to  $z_2$  becomes the small difference between the two large numbers.

This situation does not occur in practice for elements over the semi-span. However, due to the inclination of the spanwise vortex lines, it is evident from Fig. 3. that a control point could easily be aligned with the extension of an image vortex segment located in the interval  $-1 \leq z \leq 0$ . This can result in seemingly random errors which come and go with minor changes in grid spacing or planform. This problem is eliminated entirely if the velocity is set to zero for  $|d| < 0.002$  provided  $z$  is outside the interval  $z_1 < z < z_2$ . Inclusion of the small correction similar to the second term of (4.4) makes no difference, and is therefore an unnecessary complication.

The velocity induced by the two trailing vortices is:

$$\begin{aligned} \frac{4\pi v(t)(x,z)}{\Gamma} &= (z-z_1) \int_{x_1}^{\infty} \frac{d\xi}{[(x-\xi)^2 + (z-z_1)^2]^{3/2}} \\ &\quad - (z-z_2) \int_{x_2}^{\infty} \frac{d\xi}{[(x-\xi)^2 + (z-z_2)^2]^{3/2}} \\ &= \frac{1}{z-z_1} \left[ \frac{x-x_1}{\sqrt{(x-x_1)^2 + (z-z_1)^2}} + 1 \right] - \frac{1}{z-z_2} \left[ \frac{x-x_2}{\sqrt{(x-x_2)^2 + (z-z_2)^2}} + 1 \right] \end{aligned} \quad (4.5)$$

Since the trailing vortices are all parallel, the round-off error situation is much simpler than for the spanwise vortices. Difficulties could only be encountered for aspect ratios far below those of practical interest. Equation (4.5) may therefore be used for all elements.

The total velocity induced by a horseshoe element is the sum of either (4.3) or (4.4) and (4.5). The contribution of corresponding image elements

may then be obtained by repeating this computation following the substitution:

$$\begin{aligned} x_1 &\rightarrow x_2 \\ z_1 &\rightarrow -z_2 \\ x_2 &\rightarrow x_1 \\ z_2 &\rightarrow -z_1 \end{aligned} \quad (4.6)$$

which results in the correct algebraic sign for each of the individual elements. We shall use the notation  $v_{m,n,i,j}$  to denote the velocity induced at the control point located at  $(z_m, x_n)$  by the complete unit strength horseshoe element located between  $(z_i, x_{i,j})$  and  $(z_{i+1}, x_{i+1,j})$  together with its image.

The FORTRAN function HSVEL listed in the appendix performs this calculation.

The strength of the  $(i,j)$ 'th vortex element corresponding to the  $k$ 'th spanwise and  $\ell$ 'th chordwise modes of unit amplitude is

$$f_k(\tilde{z}_i) P_\ell(\tilde{s}_{i,j}), \quad (4.7)$$

where  $\tilde{z}_i$  is the transformed coordinate of the midspan of the element

$$\tilde{z}_i = \cos^{-1} \left( - \frac{(z_i + z_{i+1})}{2} \right), \quad (4.8)$$

and  $P_\ell(\tilde{s}_{i,j})$  is the integral of the chord load function obtained by substituting the  $x$ -coordinates of the leading and trailing edges of the midspan of the  $(i,j)$ 'th element as limits of integration in (3.14).

The velocity induced at the  $(m,n)$ 'th control point by the  $(k,\ell)$ 'th mode of unit amplitude is, therefore, obtained by summing the product of the mode strengths from (4.7) with the velocities for unit circulation obtained from (4.3) - (4.5) over all  $I \times J$  horseshoe elements:

$$v_{m,n,k,\ell} = \sum_{i=1}^I f_k(\hat{z}_i) \sum_{j=1}^J p_{\ell}(\hat{s}_{ij}) v_{m,n,i,j} . \quad (4.9)$$

Since the contribution of the image is included in  $v_{m,n,i,j}$ , the summation in (4.9) is only over the elements in the semi-span interval  $0 \leq z \leq 1$ .

We can now write the final expression for the velocity induced at the  $(m,n)$ 'th control point in terms of the unknown mode amplitudes  $c_{k,\ell}$  :

$$V_{m,n} = \sum_{k=1}^K \sum_{\ell=1}^L c_{k,\ell} v_{m,n,k,\ell} . \quad (4.10)$$

Equation (4.10) represents a set of simultaneous equations for the mode amplitude coefficients, once the values of  $V_{m,n}$  are prescribed by the boundary conditions of the problem. The mode amplitude coefficients corresponding to zero flap deflection and unit angle of attack,  $c_{k,\ell}^{(\alpha)}$ , are obtained by solving (4.10) with  $V_{m,n} = +1$  for all values of  $(m,n)$ . Similarly, the mode amplitude coefficients for zero angle of attack and unit flap deflection are obtained by setting  $V_{mn} = 0$  for all values of  $(m,n)$  corresponding to control points on the skeg, and  $V_{mn} = +1$  for all values of  $(m,n)$  corresponding to control points on the flap.

If the number of modes is equal to the number of control points, the boundary condition may be satisfied exactly at all the control points. If the number of modes is less than this, (4.10) may be solved by least squares to provide the closest possible fit at all the control points. The latter approach is preferred, since the higher modes amplitudes are generally of

questionable accuracy. For this application we solve (4.10) for K-L unknown mode amplitude coefficients by least squares through 80 control points.

## 5. COMPUTATION OF FORCES

Lift and induced drag may be obtained most directly from the spanwise circulation distribution given in (3.15) and (3.16), employing the well-known results of classical lifting-line theory [4]. The local lift coefficient per unit angle of attack is:

$$C_{L\alpha}(z) = \frac{-\rho U \Gamma(z)}{\frac{1}{2} \rho U^2 c(z)} = \frac{8}{c(z)} \sum_{k=1}^K f_k(z) \{c_{k1}^{(\alpha)} + c_{k2}^{(\alpha)}\}, \quad (5.1)$$

noting that the coefficients  $c_{k\ell}$  and the circulation  $\Gamma$  have been obtained from (4.10) for an angle of attack of unity. Similarly, the lift coefficient due to unit flap deflection is:

$$C_{L\delta}(z) = \frac{8}{c(z)} \sum_{k=1}^K f_k(z) \{c_{k1}^{(\delta)} + c_{k2}^{(\delta)} + c_{kL}^{(\delta)}\}. \quad (5.2)$$

Overall lift coefficients are obtained by multiplying (5.1) and (5.2) by the local chord,  $c(z)$ , integrating over the span, and dividing by the area of the lifting surface. Since only the first spanwise mode contributes, we obtain the result:

$$\begin{aligned} C_{L\alpha} &= \pi a (c_{11}^{(\alpha)} + c_{12}^{(\alpha)}) \\ C_{L\delta} &= \pi a (c_{11}^{(\delta)} + c_{12}^{(\delta)} + c_{1L}^{(\delta)}) \end{aligned}, \quad (5.3)$$

where  $a$  is the effective aspect ratio.

The induced drag coefficient, in accordance with lifting line theory, may be written as follows in terms of the present notation:



$$C_{Di\alpha} = \frac{C_{L\alpha}^2}{\pi a} \left[ 1 + \sum_{k=2}^K (2k-1) \left\{ \frac{c_{k1}^{(\alpha)} + c_{k2}^{(\alpha)}}{c_{11}^{(\alpha)} + c_{12}^{(\alpha)}} \right\}^2 \right]$$

$$C_{Di\delta} = \frac{C_{L\delta}^2}{\pi a} \left[ 1 + \sum_{k=2}^K (2k-1) \left\{ \frac{c_{k1}^{(\delta)} + c_{k2}^{(\delta)} + c_{kL}^{(\delta)}}{c_{11}^{(\delta)} + c_{12}^{(\delta)} + c_{1L}^{(\delta)}} \right\}^2 \right] \quad (5.4)$$

The efficiencies  $\eta^{\alpha, \delta}$  of the lifting-surface are defined as the reciprocals of the quantities in square brackets in (5.4), and are equal to one if the spanwise circulation coefficients are zero for  $k > 1$ .

To obtain the chordwise position of the center of pressure at any spanwise location, we must obtain the first moments of the chordwise mode functions,  $p_\ell$ .

$$M_1 = \int_{x_L}^{x_T} x \cdot p_1(\tilde{s}) \, dx = x_L(z) + \frac{c(z)}{4}$$

$$M_2 = \int_{x_L}^{x_T} x \cdot p_2(\tilde{s}) \, dx = x_L(z) + \frac{c(z)}{2}$$

$$M_3 = \int_{x_L}^{x_T} x \cdot p_3(\tilde{s}) \, dx = -\frac{c(z)}{4}$$

$$M_\ell = \int_{x_L}^x x \cdot p_\ell(\tilde{s}) \, dx = 0 \quad \text{for } \ell = 4, \dots, L-1$$

$$M_L = \int_0^{x_T} x \cdot p_L(\tilde{t}) \, dx = \frac{x_T(z)}{4} \quad (5.5)$$

The results for  $M_1$ ,  $M_2$  and  $M_L$  can be readily indentified as the positions of the center of pressure of these modes, remembering that the origin of the coordinate system is at the flap hinge, rather than at the leading edge. The third chordwise mode contributes a pure moment, while the remaining modes contribute neither force nor moment.

The local distance of the center of pressure resulting from the joint effect of all the modes is expressed as a fraction of the local chord from the hinge line as follows:

$$\frac{x_H^{(\alpha)}(z)}{c(z)} = \frac{1}{c(z)} \frac{\sum_{k=1}^K f_k(z) \sum_{\ell=1}^{L-1} c_{k,\ell}^{(\alpha)} M_\ell}{\sum_{k=1}^K f_k(z) \sum_{\ell=1}^{L-1} c_{k,\ell}^{(\alpha)} P_\ell} \quad (5.6)$$

$$\frac{x_H^{(\delta)}(z)}{c(z)} = \frac{1}{c(z)} \frac{\sum_{k=1}^K f_k(z) \sum_{\ell=1}^L c_{k,\ell}^{(\delta)} M_\ell}{\sum_{k=1}^K f_k(z) \sum_{\ell=1}^L c_{k,\ell}^{(\delta)} P_\ell} .$$

Distance between the local center of pressure and the leading edge as a function of local chord can be easily found from (5.6) as:

$$\frac{x_{LE}^{(\alpha)}(z)}{c(z)} = 1 - \frac{1}{c(z)} [x_T(z) - x_H^{(\alpha)}(z)] \quad (5.7)$$

$$\frac{x_{LE}^{(\delta)}(z)}{c(z)} = 1 - \frac{1}{c(z)} [x_T(z) - x_H^{(\delta)}(z)] .$$

Finally, the resultant chordwise position of the center of pressure relative to the flap hinge line can be obtained by integration over the semi-span,

$$\frac{\bar{x}_H^{(\alpha)}}{\bar{c}} = \frac{1}{\bar{c} C_{L\alpha}} \int_0^1 x_H^{(\alpha)}(z) C_{L\alpha}(z) c(z) dz$$

$$\frac{\bar{x}_H^{(\delta)}}{\bar{c}} = \frac{1}{\bar{c} C_{L\alpha}} \int_0^1 x_H^{(\delta)}(z) C_{L\delta}(z) c(z) dz .$$
(5.8)

This integration is performed numerically, using ten equally spaced stations and an integration formula which has been developed for functions with a square-root singularity in slope at the tip [2].

The preceding results may be combined to generate lift, drag and moment characteristics for any set of combinations of angle of attack and flap deflection. By adding an empirical viscous drag term of the form:

$$C_{DV} = 0.0085 + 0.0166 C_L^2 ,$$
(5.9)

a realistic approximation to the characteristics of a flapped control surface can be made, provided, of course, that stall does not occur. The constants appearing in (5.9) were obtained from experimental airfoil data with "standard roughness" given in [3]. A sample tabulation of this type is given in Appendix 3.

## 6. TESTS OF PROGRAM ACCURACY

The effect of the spanwise and chordwise grid spacing on computed lift-slope for a typical control surface is given in Table 3. Variations with precision numbers are quite small, and it may be concluded that zero precision numbers are satisfactory for planform shapes similar to the case examined.

Table 3

Variations of Parameters with Extreme Precision Numbers

$I_V$	$I_H$	$C_{Di\alpha}$	$C_{Di\delta}$	$C_{L\alpha}$	$C_{L\delta}$	$\eta^\alpha$	$\eta^\delta$
0	0	0.114	0.116	3.138	1.776	0.996	0.984
0	2	0.114	0.116	3.132	1.877	0.996	0.980
2	0	0.114	0.116	3.101	1.756	0.995	0.980
$a = 2.8$		$\Lambda = 15^\circ$		$f = 0.2$		$\lambda = 0.6$	

Convergence of the solution with increasing numbers of elements is a necessary but obviously not sufficient test of program accuracy. Fortunately, the results for zero flap deflection may be compared with corresponding solutions obtained by other current numerical lifting-surface theory techniques. A recent publication by Langan and Wang [5] is particularly helpful in this regard. This reference compares the spanwise distributions of lift for a tapered wing of aspect ratio 5 obtained by fifteen different lifting-surface computer programs. Fig. 5 shows the results of the present program

plotted on a reproduction of Langan and Wang's Fig. 13 which are in excellent agreement with the average results of the fifteen programs compared in [5]. The overall lift coefficient obtained by the present program for this example is 4.106 which agrees exactly with the results given in [5] obtained by the Tulinus program! Induced drag was also found to be in good agreement.

A comparison was also made of the lift coefficient of a rectangular planform of unit aspect ratio. Watkins, Woolston, and Cunningham in a 1959 report [6], cite the following values obtained from several sources, to which we have added the results from the present program:

<u>Source</u>	<u><math>C_{L\alpha}</math></u>
Jones (low aspect ratio limit)	1.571
Lawrence	1.400
Hsu	1.497
Watkins-Woolston-Cunningham	1.455
Present Program	1.508

As a further check in the low aspect ratio range, J. Dulmovits of the Grumman Aircraft Engineering Corporation kindly offered to run his program for a tapered planform of effective aspect ratio 2.8. He obtained a lift slope of 3.140, which agrees almost exactly with a value of 3.138 obtained by the present program.

A test of the portion of the program dealing with the flap was made by running a rectangular planform of aspect ratio 60, to provide an essentially two-dimensional result. The exact solution in this case, as given in [7], is

$$\frac{C_{L\delta}}{C_{L\alpha}} = \frac{[(\pi - \tilde{s}_H) + \sin(\pi - \tilde{s}_H)]}{\pi} \quad (6.1)$$

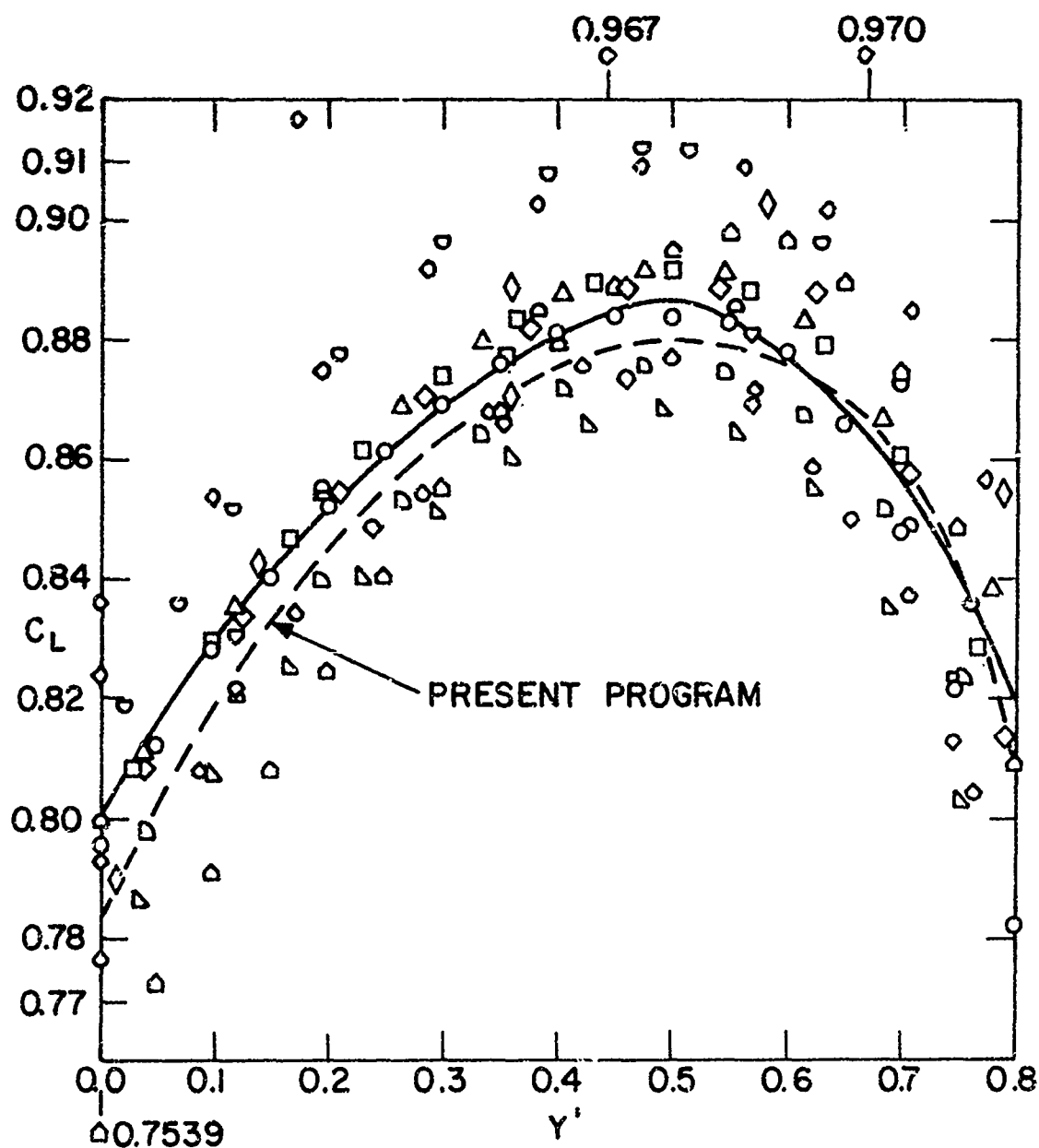


Fig. 5 Spanwise Lift Distribution Obtained by Present Program  
Superimposed on Fig. 13 of [5]

SYMBOL	PROGRAM
○	TULINUS
□	DULMOVITS
◇	MARGASON - LAMAR
△	GIESING
▤	RUBBERT
▥	LOPEZ-SHEN (EVD)
◊	HAVILAND
◻	JORDAN
◊	LAMAR (MULTHOPP)
◡	WIDNALL
◊	BANDLER (ERA)
◊	ROWE
◻	CUNNINGHAM
◊	JACOBS-TSAKONAS
◊	LOPEZ (KÜCHEMANN)

Table 4 Identification of Symbols in Fig. 5

where  $\hat{s}_H$  is the transformed chordwise coordinate evaluated at the position of the flap hinge. For a 50% flap the present program gives  $C_{L\delta}/C_{L\alpha} = .807$  compared with a value of .818 obtained from (6.1). For a 20% flap the computed value is .526 vs. .550 obtained from (6.1). We should expect slightly lower results for an aspect ratio of 60 so that this result seems reasonable.

The calculation for an aspect ratio of 60 also provided an additional check on  $C_{L\alpha}$  in the limit of high aspect ratio. The value obtained was 5.915 which compares closely with a value of 6.080 which one would obtain from the lifting-line theory.

It may also be of some interest to compare the value of  $C_{L\alpha}$  given by the present program with that given by Equation (23a) of Chapter VIII of Principles of Naval Architecture [8]. The latter is a simple empirical equation:

$$C_{L\alpha} = \frac{(0.9) (2\pi)a}{(\cos\Lambda \sqrt{\frac{a^2}{\cos^4\Lambda} + 4}) + 1.8} \quad (6.2)$$

which, for an aspect ratio  $a = 2.8$  and a sweep angle  $\Lambda = 15^\circ$  gives a value of 2.997, which is about five percent below the computed value of 3.138. Since (6.2) was developed to provide correlation with experimental results, the lower value of lift slope is not at all unreasonable.



## 7. DESIGN APPLICATION

The present program was initially developed to provide a rational basis for selecting optimum planforms for an experimental series of flapped rudders [9]. The first two rudders in this series were to have an effective aspect ratio of 2.8 and flap area ratios of 20% and 10%, respectively. These were planned as an extension of an earlier series of experiments on flapped rudders published in 1972 [10].

With such small flaps, minor changes in sweep angle and taper ratio have a large effect on the spanwise distribution of flap chord. The objective is to determine the sweep angle and taper ratio in such a way as to optimise the following parameters:

- a) Maximize lift slopes  $C_{L\alpha}$  and  $C_{L\delta}$  ;
- b) Minimize induced drag;
- c) Provide nearly uniform distribution of  $C_L$  over span.  
This requirement is based on a decision to use a constant thickness/chord ratio of 15% over the span. Uniform lift coefficient is then optimum both for delay of cavitation inception and delay of stall.
- d) Provide sufficient flap tip chord to permit installation of a hinge.

Calculations were made for five combinations of taper ratio and sweep for the 20% flap rudder and the principal results are given in Table 5. Plots of spanwise distribution of lift-slopes are shown in Fig. 6.

Table 5

Effect of Sweep and Taper on Flapped Rudder Characteristics  
effective aspect ratio = 2.8

No.	f	$\Lambda$	$\lambda$	$C_{L\alpha}$	$C_{L\delta}$	$C_{D\alpha}$	$C_{D\delta}$
1	0.2	11°	0.9	3.101	1.670	0.115	0.121
2	"	11°	0.6	3.146	1.802	0.114	0.115
3	"	15°	0.6	3.138	1.775	0.114	0.116
4	"	18°	0.6	3.129	1.738	0.114	0.117
5	"	19.57°	0.5	3.136	1.774	0.114	0.115
6	0.1	15°	0.6	3.138	1.355	0.114	0.119

It is clear from Table 5 that overall lift and drag characteristics are very insensitive to sweep and taper within the fairly limited range permitted due to the small flap. This is, of course, a characteristic of low aspect ratio lifting surfaces, so that this result is not surprising. If one looks closely, one can see the trend of decreasing lift slope and increasing drag with increasing sweep angle.

The effect of taper ratio on spanwise distribution of lift is much more pronounced, as is evident from Fig. 6. Here it is clear that a taper ratio of 0.9 unloads the tip too much, while a taper ratio of 0.5 does the reverse, and that a value of 0.6 seems about right. A change in sweep angle from 11 to 18 degrees has essentially no effect on  $C_{L\alpha}$ , and a small effect on  $C_{L\delta}$ .

Rudder No. 2 appears to have the best spanwise distribution of lift as well as maximum lift and minimum drag. However, the tip chord of the flap is extremely small, which would cause difficulties in the model and possibly

in the full-scale hinge design. An increase in sweep angle from 11 to 15 degrees overcomes this problem with very little compromise in performance. Rudder No. 3 was therefore selected for the test program. This has the additional practical advantage of a nearly vertical trailing edge and a constant flap chord. .

The constraint on flap chord for the 10% flap area ratio rudder is even more severe. Rudder No. 6 shown in Table 5 and Fig. 6 has the same planform as No. 3, but with a flap area of 10%. Again, we find that  $C_{L\delta}$  is a little too low at  $z = 0$ . Increasing this by a decrease in sweep angle is even more impractical in this case due to the small flap chord. We conclude, therefore, that this planform seems to be nearly optimum for both the 20% and 10% flap rudders.

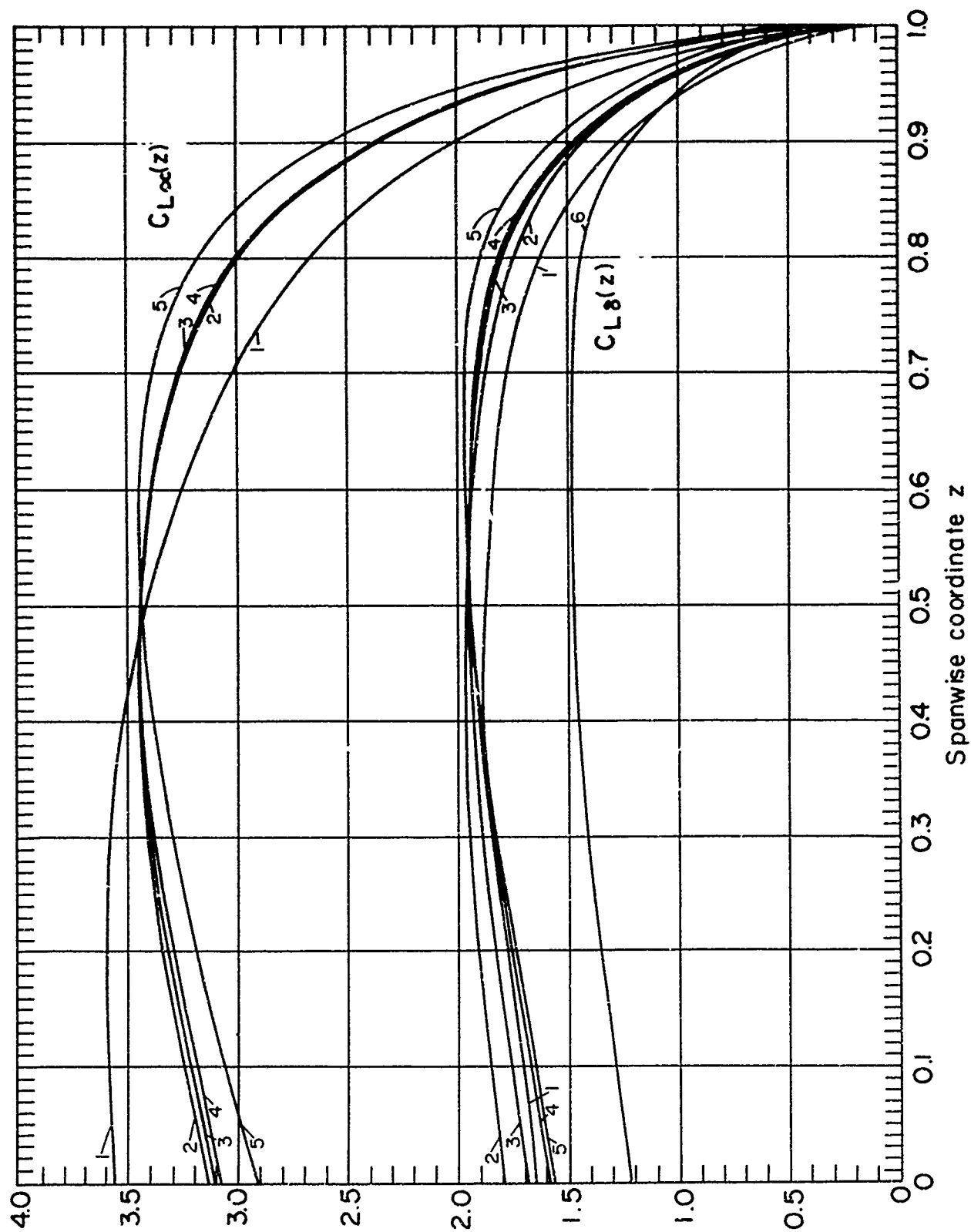


Fig. 6 Spanwise Distribution of Lift due to Angle of Attack and Flap Deflection for Rudder Designs of Table 6

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A P P E N D I C E S

## APPENDIX 1

### Instructions for Preparing Computer Program Input Data

Integer variables must be right-adjusted in their fields - Real variables must contain a decimal somewhere within their field.

#### A. FIRST CARD

SYMBOL	MODE	FIELD	LIMITATIONS	DESCRIPTION
KDM	Integer	4	1-6	Number of spanwise modes-K Recommended value is 6
LT	Integer	8	3	Number of chordwise modes-L Recommended value is 6
IHF	Integer	12	0,1 or 2	Chordwise precision index $I_H$ Recommended value is 0
IV	Integer	16	0,1 or 2	Spanwise precision index $I_V$ Recommended value is 0
ASR	Real	17-24	*	Geometric aspect ratio, $\frac{1}{2}a$
AF	Real	25-32	$0.1 \leq AF \leq 0.9$	Flap area ratio, f
T	Real	33-40	*	Taper ratio, $\lambda$
PP	Real	41-48	*	Sweep angle in degrees, $\Lambda$
KOPT	Integer	56	0 or 1	If KOPT=1, the subroutine OPTION will be called and will perform the calculations appearing on the third page of the output

#### B. SECOND CARD

This card contains the indices of the ten chordwise panels which contain control points, NCP(N), N=1,10. These are integers which are read in ten

\*Combinations of these quantities must be so chosen that flap chords are positive for all z.

fields of 8 columns. The recommended values for zero precision number are 3, 8, 13, 18, 23, 28, 33, 38, 42, and 48. However, the intention is that these may be altered for a more advantageous placement relative to the flap hinge if so desired.

C. MULTIPLE RUNS

Upon completion of the calculation, the program returns to the read statement for the first card. A blank card (or more specifically, a zero value for KDM) terminates the run.



# APPENDIX 2

## Sample Output

The correspondence between the output and the text symbols is as follows:

text	output	text	output
$z$	$\bar{z}$	$C_{M\alpha}(z) = \frac{x_H^{(\alpha)}(z)}{c(z)} \cdot C_{L\alpha}(z)$	CM-A
$C_{L\alpha}(z)$	CLA	$C_{M\delta}(z) = \frac{x_H^{(\delta)}(z)}{c(z)} \cdot C_{L\delta}(z)$	CM-D
$C_{L\delta}(z)$	CLD	$\alpha$	ALPHA
$C_{L\alpha}$	CLAR	$\delta$	DELTA
$C_{L\delta}$	CLDR	$\delta/\alpha$	DELTA/ALPHA
$C_{Di\alpha}/C_{L\alpha}^2$	CDIA	$C_L = C_{L\alpha} \cdot \alpha + C_{L\delta} \cdot \delta$	CL
$C_{Di\delta}/C_{L\delta}^2$	CDID	$C_{Di} + C_{DV}$	CD
$\eta^\alpha$	EFA	$(C_{L\alpha} \cdot C_M \cdot \alpha^2 + C_{L\delta} \cdot C_{L\delta} \cdot \delta^2)/C_L$	CM
$\eta^\delta$	EFD	$\frac{x_H^\alpha}{c} \cdot C_{L\alpha} \cdot \alpha + \frac{x_H^\delta}{c} \cdot C_{L\delta} \cdot \delta)/C_L$	XHL/C
$V_{m,n}$	downwash velocities		
$c_{KL}^\alpha$	C-ALPHA		
$c_{KL}^\delta$	C-DELTA		
$\frac{x_H^{(\alpha)}(z)}{c(z)}$	XA/LC		
$\frac{x_H^{(\delta)}(z)}{c(z)}$	XD/LC		
$\frac{x_{LE}^{(\alpha)}(z)}{c(z)}$	XALE/LC		
$\frac{x_{LE}^{(\delta)}(z)}{c(z)}$	XDLE/LC		

ASPECT RATIO= 2.80 FLAP AREA=200 TAPER RATIO=0.60 SHEEP ANGLE=15.000 DEG.  
 PRECISION INDICES: CHORDWISE = 0 SPANWISE = 0  
 NUMBER OF SPANWISE MODES = 6 NUMBER OF CHORDWISE MODES = 6

# SPANWISE DISTRIBUTION OF LIFT

Z = SPAN COORDINATE (Z=0 AT THE POINT, Z=1 AT THE TIP)

Z	CLA	CLC
0.0	3.104	1.693
0.050	3.164	1.717
0.100	3.220	1.753
0.150	3.271	1.789
0.200	3.317	1.824
0.250	3.357	1.856
0.300	3.389	1.884
0.350	3.413	1.907
0.400	3.428	1.923
0.450	3.434	1.931
0.500	3.429	1.937
0.550	3.412	1.935
0.600	3.382	1.929
0.650	3.335	1.916
0.700	3.264	1.894
0.750	3.159	1.854
0.800	3.001	1.783
0.850	2.758	1.661
0.900	2.380	1.456
0.925	2.111	1.304
0.950	1.758	1.100
0.975	1.262	0.803
0.995	0.568	0.368
1.000	0.0	0.0

OVERALL LIFT SLOPE COEFF. PER RADIAN CLAP= 3.138 CLOR= 1.776  
 INDUCED DRAG COEFF./UNIT LIFT COEFF.=0.2 CLTA= 0.114 CIND= 0.114  
 EFFICIENCIES ARE RESPECTIVELY LFA= 0.298 EFD= 0.984

## K DENOTES SPANWISE MODE, L DENOTES CHORDWISE MODE

K DEVIATIONS SPANWISE		L DEVIATIONS CHORDWISE					
MODE		MODE					
K=1	K=2	K=3	K=4	K=5	K=6		
L=1	0.38337	0.05644	0.01467	0.01509	0.00131	0.00143	
L=2	-0.02663	-0.04391	-0.01707	-0.01471	-0.00242	-0.00148	
L=3	0.00410	-0.01519	-0.01336	-0.01192	-0.00436	-0.00260	
L=4	0.00493	-0.00349	-0.00435	-0.00541	-0.00315	-0.00154	
L=5	0.00344	-0.00022	-0.00254	-0.00274	-0.00237	-0.00124	
NONE		C-DELTA					
K=1	K=2	K=3	K=4	K=5	K=6		
L=1	0.03184	0.00799	0.00094	-0.00145	0.00271	-0.00050	
L=2	0.12911	0.00303	-0.00524	-0.00001	-0.00589	0.00109	
L=3	-0.02936	0.00955	0.00154	0.00312	-0.00229	-0.00027	
L=4	0.03670	0.00281	-0.00339	0.00453	-0.00563	0.00244	
L=5	0.02933	0.00497	0.00748	0.00017	-0.00147	-0.00109	
L=6	0.04197	0.00845	0.00605	0.00140	0.00213	-0.00024	

ASPECT RATIO=2.80 FLAP AREA=200 TAPER RATIO=0.60 SWEEP ANGLE=15.000 DEG.  
 PRECISION INDICES: CHORDWISE=0 SPANWISE=0  
 NUMBER OF SPANWISE MODES = 6 NUMBER OF CHORDWISE MODES = 6

MATRIX OF DOWNWASH VELOCITIES AT ALL CONTROL POINTS  
 M DENOTES SPANWISE INDEX, N DENOTES CHORDWISE INDEX  
 CUE TO ALPHA

	N=1	N=2	N=3	N=4	N=5	N=6	N=7	N=8	N=9	N=10
M=1	1.02080	1.01116	1.01427	1.00965	1.00105	0.99899	0.99847	0.99976	1.00008	0.99609
M=2	0.96733	0.96239	0.97996	0.98733	0.99302	0.99241	0.99632	1.00106	1.00412	1.00172
M=3	1.02162	1.00903	1.01750	1.01494	1.00614	1.00314	1.00235	1.00461	1.00584	0.99946
M=4	1.00681	0.99565	1.00701	1.00433	0.99656	0.99097	0.99077	0.99471	0.99770	0.99034
M=5	0.99270	0.97749	0.99474	0.99995	0.99997	0.99854	1.00146	1.00599	1.01031	1.00458
M=6	1.01415	0.99568	1.01201	1.01273	1.00253	0.99502	0.99274	0.99644	1.00129	0.99122
M=7	1.00180	0.97540	1.00181	1.00650	1.00174	0.99539	0.99479	1.00150	1.01019	0.99965
M=8	0.99924	1.00442	1.00358	0.99766	0.99761	1.00075	1.00275	1.00064	0.99525	1.00047
CUE TO DELTA										
	N=1	N=2	N=3	N=4	N=5	N=6	N=7	N=8	N=9	N=10
M=1	-0.03657	0.06400	0.02941	-0.00950	-0.02793	-0.00497	0.05994	-0.01547	0.93995	1.03361
M=2	-0.01309	0.09795	-0.03209	-0.12578	-0.06381	0.06318	0.15202	-0.01752	0.87755	1.05137
M=3	-0.06584	0.12095	0.01788	-0.07387	-0.05676	0.01937	0.12188	-0.01365	0.85965	1.07656
M=4	-0.05024	0.12134	0.02771	-0.02038	-0.06954	0.01893	0.13329	-0.01200	0.84314	1.05579
M=5	-0.03746	0.11904	-0.01552	-0.11144	-0.08395	0.04464	0.15288	-0.00392	0.83164	1.05077
M=6	-0.05973	0.12796	0.01428	-0.09097	-0.07694	0.01438	0.14183	0.00333	0.86791	1.04599
M=7	-0.07088	0.12723	0.01498	-0.09466	-0.08966	0.02297	0.13937	0.00422	0.89483	1.04081
M=8	-0.05432	0.12162	0.01434	-0.09399	-0.08150	0.01343	0.13969	-0.00431	0.88647	1.04294

CONTROL POINT COLUMNS ARE LOCATED AT THE DOWNSTREAM BOUNDARIES OF THE FOLLOWING PANELS

3	9	13	18	23	24	33	38	43	48
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ASPECT RATIO= 2.80 FLAP AREA=.200 TAPER RATIO=0.60 SWEEP ANGLE=15.000DEG. PRECISION INDICES: CHORDWISE =0 SPANWISE=0  
 NUMBER OF SPANWISE MODES = 6 NUMBER OF CHORDWISE MODES = 6

ALPHA	DELTA	DFL/ALPHA	CL	CD	L/D	CM	XHL/C
0.0	0.0	0.0	0.0	0.0095	0.0	0.0	0.00000
0.0	5.0	0.0000	0.1549	0.0117	13.275	-0.07627	-0.2341
0.0	10.0	0.0000	0.3099	0.0212	15.627	-0.07754	-0.2341
0.0	15.0	0.0000	0.4648	0.0370	15.548	-0.10880	-0.2341
0.0	20.0	0.0000	0.6198	0.0597	10.461	-0.14507	-0.2341
0.0	25.0	0.0000	0.7747	0.0878	8.825	-0.18134	-0.2341
5.0	0.0	0.0	0.2718	0.0183	14.961	-0.15737	-0.5747
5.0	2.5	0.5	0.3513	0.0247	15.744	-0.12667	-0.4996
5.0	5.0	1.0	0.4288	0.0326	13.150	-0.11161	-0.4516
5.0	7.5	1.5	0.5063	0.0421	12.014	-0.11010	-0.4183
5.0	10.0	2.0	0.5837	0.0533	10.761	-0.11233	-0.3939
5.0	12.5	2.5	0.6612	0.0660	10.025	-0.11829	-0.3751
10.0	0.0	0.0	0.7477	0.0477	11.478	-0.31474	-0.5747
10.0	5.0	0.5	0.7026	0.0732	9.605	-0.25333	-0.4996
10.0	10.0	1.0	0.6576	0.1049	8.173	-0.22720	-0.4516
10.0	15.0	1.5	1.0125	0.1431	7.078	-0.22320	-0.4183
10.0	20.0	2.0	1.1675	0.1875	6.226	-0.22467	-0.3939
10.0	25.0	2.5	1.3224	0.2383	5.549	-0.21659	-0.3751
15.0	0.0	0.0	0.8215	0.0967	8.492	-0.47211	-0.5747
15.0	7.5	0.5	1.0540	0.1540	6.845	-0.38000	-0.4996
15.0	15.0	1.0	1.2464	0.2255	5.705	-0.34083	-0.4516
15.0	22.5	1.5	1.5188	0.3112	4.880	-0.31030	-0.4183
15.0	30.0	2.0	1.7512	0.4113	4.258	-0.33700	-0.3939
15.0	37.5	2.5	1.9836	0.5256	3.774	-0.35489	-0.3751
20.0	0.0	0.0	1.0954	0.1654	6.624	-0.62948	-0.5747
20.0	10.0	0.5	1.4053	0.2671	5.261	-0.50667	-0.4996
20.0	20.0	1.0	1.7152	0.3542	4.351	-0.45444	-0.4516
20.0	30.0	1.5	2.0750	0.5467	3.704	-0.44040	-0.4183
20.0	40.0	2.0	2.3349	0.7246	3.222	-0.44934	-0.3939

INTEGRATED

SPANWISE POSITIONS OF THE CENTERS OF PRESSURE ARE : ZA= 0.44 ZD= 0.47

### APPENDIX 3

#### Computer Program Listing and Particulars

Timing Information: 0 precision with 36 modes;  
1.081 minutes execution time.

IBM System 370/165

Memory Requirement 100K

47

```

0050      NS1=NS+1
0051      NT=NS+NF
0052
0053      17      XCP(I,J)=XF*(FLOAT(J-NS-1)*0.90)/FLOAT(NF)
0054      XCP=ZCP(K)*STR+SRGT
0055      XFCP=ZCP(K)*FTR+FRONT
0056      ON 20 N=1,10
0057      IF(NCP(N).GT.NS) GO TO 19
0058      XCP(K,N)=XFCP+FLOAT(NCP(N)-NS)/FLOAT(NS)
0059      XCP(K,N)=XSCP+FLOAT(NS-NCP(N))/FLOAT(NF)
0060      GO TO 20
0061      19      XCP(K,N)=XFCP+FLOAT(NCP(N)-NS)/FLOAT(NS)
0062      GO TO 20
0063      C      CALCULATION OF INDUCED VFLOCITIES *****
0064      GO 1100 KD=1,NUNK
0065      CA(KD)=0.0
0066      CD(KD)=0.0
0067      DD 759 KD=1,KDM
0068      DD 759 I=1,LT
0069      DD 759 K=1,9
0070      DD 759 N=1,10
0071      VIF,A,KD,L)=0.0
0072      LWS=LX2-1
0073      ZWS=0.5*(ZCP(I+1)+ZCP(I))
0074      XS=ZMS*STR+SRGT
0075      XF=ZMS*FTR+FRONT
0076      CGPT=ZMS*(FTR-STR)*FRGCT-SRGT
0077      ZWIG=ARCOS(1-ZWS)
0078      ON 500 K9=1,KM
0079      E(K9)=SIN(12.0*(FLOAT(KD)-1.0)*ZWIG)
0080      DD 720 J=1,NT
0081      FJ=FLFAT(J)
0082      DD 721 K=1,9
0083      DD 721 N=1,10
0084      ICM=N1
0085      JCM=N2
0086      KCM=N3
0087      NCM=N4
0088      721      Q(K,N)=HVEL(XCP(I,J),ZCP(I),XCP(I+1),XCP(K,N),ZCP(K))
0089      IF (J.GT.NS) GO TO 600
0090      SS2=XSCFJ/INSCGPT
0091      SS1=XSCFJ*(1.0)/INSCGPT
0092      SSW=ARCOS(1.0-2.0*SS2)
0093      SSW1=ARCOS(1.0-2.0*SS1)
0094      P(1)=SSW2+SIN(SSW2)-SSW1-SIN(SSW1)
0095      P(2)=SSW2-0.5*SIN(2.0*SSW2)-SSW1+0.5*SIN(2.0*SSW1)
0096      IF (1.1E-3) GO TO 710
0097      DD 605 I=3,11
0098      FL=FLOAT(I)
0099      609      P(1)=SIN(FL-2.0)*SSW2-SIN(FL-2.0)*SSW1)/((FL-2.0)*(SIN(FL+SSW1)
0100      1)-SIN(FL+SSW2))/FI
      GO TO 710

```

```

0101 600 FJOT=FLOAT(J-NS)
0102 SF2=(XF*FJOT-RNF*XS)/(RNF*CGPT)
0103 SF1=(XF*(FJOT-1.0)-RNF*XS)/(RNF*CGPT)
0104 SFW2=ARCS(1.0-2.0*SF2)
0105 SFW1=ARCS(1.0-2.0*SF1)
0106 P(1)=SFW2+SIN(SFW2)-SFW1-SIN(SFW1)
0107 P(2)=SFW2-C.5*SIN(2.0*SFW2)-SFW1*0.5*SIN(2.0*SFW1)
0108 IF (17.4E-3) GO TO 700
0109 DO 704 L=3,LL
0110 FL=FLOAT(1)
0111 P(1)=SIN((FL-2.0)*SFW2)-SIN((FL-2.0)*SFW1)/(FL-2.0)+(SIN(FL*SFW1)
0112 )-SIN(FL*SFW2))/FL
0113 T2=FJOT/RNF
0114 T1=(FJOT-1.0)/RNF
0115 TW2=ARCS(1.0-2.0*T2)
0116 TW1=ARCS(1.0-2.0*T1)
0117 LIT=1
0118 P(LT)=TW2+SIN(TW2)-TW1-SIN(TW1)
0119 DO 760 K=1,8
0120 DO 760 N=1,10
0121 DO 760 KD=1,KDM
0122 V(K,N,KD,L)=V(K,N,KD,L)+F(KD,N)*F(KD)*P(L)
0123 CONTINUE
0124 C CALCULATION OF BOUNDARY CONDITIONS *****
0125 DO 801 K=1,9
0126 DO 801 N=1,10
0127 DNMWSA(K,N)=0.0
0128 DNMWS(K,N)=0.0
0129 LTA=LT-1
0130 NUNK=KDM*(LT-1)
0131 MUNK=NUNK+1
0132 DO 805 N=1,10
0133 IV=(N-1)*P+K
0134 DO 800 L=1,LT5
0135 DO 800 KD=1,KDM
0136 IV=(L-1)*KD+KD
0137 W(I,V,IV)=V(K,N,KD,L)
0138 W(I,V,MUNK)=1.0
0139 CALL PTLQ (W,CA,80,MUNK,ERROR)
0140 DO 803 N=1,10
0141 DO 803 K=1,8
0142 DO 803 KD=1,KDM
0143 DO 803 L=1,LT5
0144 IS=(L-1)*KD+KD
0145 DNMWSA(K,N)=V(K,N,KD,L)*CA(15)+DNMWSA(K,N)
0146 NUNK=KD+LT
0147 MUNK=NUNK+1
0148 DO 815 N=1,10
0149 DO 815 K=1,8
0150 IV=(N-1)*K+K
0151 DO 810 L=1,LT

```



```

0152 DD 910 KD=1,KDM
0153 IH=(L-1)*KDM*KD
0154 MW(IV,IH)=V(K,KD,L)
0155 W(IIV,MUNK)=1,C
0156 IFIXCPIK(A).LT.0.0) MW(IIV,MUNK)=0.0
0157 CONTINUE
0158 CALL PYLSQ(MW,FC,RO,MUNK,KFRRO)
0159 DO 904 K=1,10
0160 DO 904 K=1,R
0161 DO 904 KD=1,KDM
0162 DO 904 L=1,LT
0163 DO 904 L=1,LT
0164 DO 904 L=1,LT
0165 DNMWSDIK(N)=V(K,KD,L)*C*(IS)+DNMWSDIK(N)
0166 C CALCULATIONS OF FORCES *****
0167 CLG=(3.14159*3)*AR*(C(I)*C(I)+C(I)*C(I)+C(I)*C(I)+C(I)*C(I))
0168 SM1=0.0
0169 SM2=0.0
0170 ZMP(I)=0.0
0171 DO 910 I=2,77
0172 IF (I-LE-1) ZMP(I)=ZMP(I-1)+0.05
0173 IF (I-CT-1) ZMP(I)=ZMP(I-1)+0.075
0174 CONTINUE
0175 ZMP(23)=0.995
0176 ZMP(24)=1.C
0177 CLA(24)=0.0
0178 CLD(24)=0.0
0179 DO 911 I=1,23
0180 SUM1=0.0
0181 SUM2=0.0
0182 CPT=ZMP(I)*(FIR-STR)*FRCCT-SROOT
0183 ZWIG=APCOS(-ZMP(I))
0184 DO 912 KD=1,KDM
0185 FD=F(CAT(KC))
0186 EIKD)=STN(I2.0*FK-1.0)*ZWIG)
0187 SUM1=SUM1+E(IK)*CA(KD)*CA(KD+KDM)*2.0/CGPT
0188 SUM2=SUM2+E(IK)*CA(KD)*CA(KD+KDM)*C(I)*C(I)*NDX)*2.0/CGPT
0189 CLA(I)=39.47842*SUM1
0190 CLD(I)=39.47842*SUM2
0191 AA1=CA(I)*CA(I)*KDM
0192 AA2=CD(I)*C(I)*KDM+CD(I)*NDX
0193 DO 913 KD=2,KDM
0194 FK=FINAT(KC)
0195 SM1=SM1+(2.0*FK-1.0)*((CA(KD)+CA(KD+KDM))/AA1)**2
0196 SM2=SM2+(2.0*FK-1.0)*((CD(KD)+CD(KD+KDM))/AA2)**2
0197 COTA=(1.0+SM1)/(1.4159*AR)
0198 CDIN=(1.0+SM2)/(1.4159*AR)
0199 EFD=1.0/(1.0+SM2)
0200 KD=1,KDM*LT
0201 DO 915 N=1,KC*L
0202 CAA(N)=CA(N)*9.86958
0203 CND(N)=CND(N)*9.86958
015 CND(N)=CND(N)*9.86958

```

```

0204 WRITE(6,1003)AP,AF,T,PP,IMF,IVER,KDP,LT
0205 1-03 FORMAT(11,5X,17HASPECT RATIO=FS,2,2X,10HFLAP AREA=FA,3,2X,12HTAPE
1R RATIO=FA,2,2X,12HSWEEP ANGLE=FA,3,3X,4HDEG.,7,6X,20HPERFECTSION IN
2DICES:CHORDWISE =11,2X,9HSPANWISE=11,7,6X,10HNUMBER OF SPANWISE MD
3DES =12,5X,10HNUMBER OF CHORDWISE WINGS =12,77)
0206 WRITE (6,548)
0207 FORMAT (10X,20HSPANWISE DISTRIBUTION OF LIFT,7,10X,2 = SPAN COORD
1INATE (2=0 AT THE ROOT,2 =1 AT THE TIP),777)
0209 WRITE(6,946)
0210 FORMAT(16X,1M7,14X,3HCLA,12X,3HCLD)
0211 DC 950 1=1,24
0212 WRITE(6,945)7=0(1),CLA(1),CLD(1)
0213 FORMAT (10X,3(F10.3,5X))
0214 WRITE(6,991)CLAG,CLDG
0215 FORMAT(7,9X,10HVEFALL LIFT SLOPE COEFF. PER RADIAN CLAR=,F0.3,
16X,1CLCP=,F8.3)
0216 WRITE(6,1000)CIN,CIDN
1000 FORMAT(9X,10HINDUCED DRAG COEFF./UNIT LIFT COEFF.,002 CDA=,F8.3,6X,
11C)ID=,F8.3)
0217 WRITE(6,961)FEA,FFD
0218 FORMAT(9X,20HEFFICIENCIES ARE RESPECTIVELY,12X,4HEFA=FA,3,6X,5HEFD
1=F8.3)
0219 WRITE(6,955)
0220 FORMAT (7,10X,10HDF AMPLITUDES C-ALPHA=)
0221 WRITE(6,1002)
1002 FORMAT(9X,50HX DENOTES SPANWISE MODE, L DENOTES CHORDWISE MODE)
0223 WRITE(6,1001)
1001 FORMAT (18X,3HK=1,9X,3HK=2,9X,3HK=3,9X,3HK=4,9X,3HK=5,9X,3HK=6)
0225 DO 957 L=1,15
0226 KX=(L-1)*KDM+1
0227 KMAX=KX+KCM-1
0228 FORMAT (9X,2HL,11,1X,10F12.5)
0229 WRITE(6,958)1,11,1X,10F12.5)
0230 WRITE(6,959)
0231 FORMAT (7,10X,10HDF AMPLITUDES C-DELTA=)
0232 WRITE(6,1001)
0233 DO 960 L=1,LT
0234 KX=(L-1)*KDM+1
0235 KMAX=KX+KCM-1
0236 WRITE(6,958)1,11,1X,10F12.5)
0237 WRITE(6,1001) AP,AF,T,PP,IMF,IVER,KDP,LT
0238 WRITE (6,541)
0239 FORMAT (7,10X,10HDF AMPLITUDES C-DELTA=)
0240 15,7,20X,10HDF AMPLITUDES C-DELTA=,11,1X,10F12.5)
0241 230X,10HDF TO ALPHA
0242 3,9X,1N=5,9X,1N=6,9X,1N=7,9X,1N=8,9X,1N=9,9X,1N=10)
0243 DO 982 K=1,8
0244 WRITE(6,983)K,(CMN$A(K),N=1,10)
0245 982 WRITE(6,983)K,(CMN$A(K),N=1,10)
0246 983 FORMAT (9X,1N=1,1X,10F12.5)
0247 WRITE(6,984)
0248 984 FORMAT (7,10X,10HDF TO DELTA=,11,1X,10F12.5)
0249 1,9X,1N=5,9X,1N=6,9X,1N=7,9X,1N=8,9X,1N=9,9X,1N=10)
0250 DO 985 K=1,8

```

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```
FORTRAN IV G1 RELEASE 2.0      MAIN      DATE = 74226      18/06/24
0246      925 WRITE (6,925) (CMNNSP(K,M),M=1,10)
0247      WRITE(6,926) (ICP(N),N=1,10)
0248      FORMAT (////,10X,'CONTROL POINT COLUMNS ARE LOCATED AT THE NOMNST
0249      IREAN ROUNDCARTS OF THE FOLLOWING PANELS :',10X,10I10,///)
0250      IF(KOPT.EQ.1) CALL DITION (CA,CD,CM,LT,INF,IVER,AR,AP,T,PP,SRONT,
0251      LSTIP,FRONT,FTIP)
      GO TO 2
      END
```

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PTLSQ

```
0001 SURRCUTINE PTLSQ (A,R,NEC,MIV,KERROR)
0002 DIMENSION A(80,37),R(36),B(144)
0003 MUK=MUK+1
0004 DO 1 M=1,NUN
0005 DO 1 A=1,NUN
0006 L=M*(M-1)*NUN
0007 R(L)=0.0
0008 DO 1 J=1,NEO
0009 B(L)=R(L)+A(J,M)*A(J,N)
0010 DO 2 M=1,NUN
0011 R(M)=0.0
0012 DO 2 N=1,NEO
0013 R(M)=R(M)+A(A,MIN)*A(N,M)
0014 CALL SLNG (R,P,MUN,KERROR)
0015 RETURN
0016 END
```

1

2

```

0001 SUBROUTINE SIMO(A,N,KSI)
0002 DIMENSION A(1),B(1)
0003 TOL=0.0
0004 KS=0
0005 JJ=N
0006 DO 65 J=1,N
0007 JY=J+1
0008 JJ=JJ+N+1
0009 AIGA=0
0010 IT=JJ-J
0011 DO 10 I=J,N
0012 IJ=IT+I
0013 IF(ABS(B(IJA))-ABS(A(IJ)))20,30,30
0014 20 AIGA=A(IJ)
0015 IMAX=I
0016 30 CONTINUE
0017 IF(IARS(IIGA)-TCI)25,35,40
0018 35 KS=1
0019 RETURN
0020 40 I1=J+N*(J-2)
0021 IT=IMAX-J
0022 DO 50 K=J,N
0023 I1=I1+N
0024 I2=I1+IT
0025 SAVE=A(I1)
0026 A(I1)=A(I2)
0027 A(I2)=SAVE
0028 50 A(I1)=A(I1)/AIGA
0029 SAVE=A(I1)
0030 R(I1)=-9(J)
0031 R(J)=SAVE/AICA
0032 IF(J=N) 55,70,55
0033 55 IQS=N*(J-1)
0034 DO 65 IX=JY,N
0035 IJ=IQS+IX
0036 IT=J-IX
0037 DO 60 JX=JY,N
0038 IJX=N*(JX-1)+IX
0039 JJX=IJX+IT
0040 60 A(IJX)=A(IJX)-(A(IJX)*A(IJX))
0041 65 R(IJX)=R(IJX)-(R(IJ)*A(IJX))
0042 70 NY=N-1
0043 IT=N*N
0044 DO 80 J=1,NY
0045 IA=IT-J
0046 IR=N-J
0047 IC=N
0048 DO 80 K=1,J
0049 R(IA)=R(IR)-A(IA)*R(IC)
0050 IA=IA-N
0051 IC=IC-1
0052 RETURN
0053 END

```

SIMO 490  
 SIMO 500  
 SIMO 540  
 SIMO 550  
 SIMO 560  
 SIMO 570  
 SIMO 580  
 SIMO 590  
 SIMO 600  
 SIMO 610  
 SIMO 620  
 SIMO 660  
  
 SIMO 680  
 SIMO 690  
 SIMO 700  
  
 SIMO 750  
 SIMO 760  
 SIMO 800  
 SIMO 810  
 SIMO 820  
 SIMO 830  
 SIMO 840  
 SIMO 850  
 SIMO 860  
 SIMO 870  
 SIMO 910  
 SIMO 920  
 SIMO 930  
 SIMO 940  
 SIMO 980  
 SIMO 990  
 SIMO1000  
 SIMO1010  
 SIMO1020  
 SIMO1030  
 SIMO1040  
 SIMO1050  
 SIMO1060  
 SIMO1070  
 SIMO110  
 SIMO1120  
 SIMO1130  
 SIMO1140  
 SIMO1150  
 SIMO1160  
 SIMO1170  
 SIMO1180  
 SIMO1190  
 SIMO1200  
 SIMO1210  
 SIMO1220

```

FUNCTION HVEL(X1,Z1,X2,Z2,X,Z)
HVEL=0.0
XA=X1
ZA=Z1
XB=X2
ZB=Z2
DO 1 N=1,2
T=(XR-XA)/(ZB-ZA)
A=1.0+T**2
E=X-XA+T*ZA
B=-2.0*(E*T+Z)
C=E**2+Z**2
D=E-T*Z
RAT=1.0
ZC=ZA
XC=XA
DO 2 M=1,2
HVEL=HVEL+RAT*((X-XC)/SQRT((X-XC)**2+(Z-ZC)**2)+1.0)/(Z-ZC)
IF(ABS(D).GT.0.002) GO TO 3
IF(N.NE.1.OR.Z.LT.Z1.OR.Z.GT.Z2) GO TO 4
HVEL=HVEL+SQRT(A)*(1.0/D-4.0*D/(2.0*A*ZC+B)**2)
GO TO 4
3 HVEL=HVEL-RAT*0.5*(2.0*A*ZC+B)/(D*SQRT(A*ZC**2+B*ZC+C))
4 RAT=-1.0
XC=XB
ZC=ZB
XA=X2
ZA=Z2
XB=X1
ZB=Z1
RETURN
END

```



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10/06/74

DATE = 74226

OPT(CM

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```
0048 DO 1525 L=4,LW8
0049 CM(L)=0.0
0050 CM(LT)=3.14159*XT/4.0
0051 CMAL=0.0
0052 CML=0.0
0053 CHA=0.0
0054 CHD=0.0
0055 DO 1540 KD=1,KDM
0056 CMAK=0.0
0057 CMDK=0.0
0058 CHAK=0.0
0059 CHDK=0.0
0060 DO 1530 L=1,LW8
0061 MP=PD+KD*(L-1)
0062 CHAK=CHAK+CA(MP)*D(L)
0063 CMAK=CMAK+CM(L)*CA(MP)
0064 DO 1535 L=1,LT
0065 MP=KD+KDM*(L-1)
0066 CHDK=CHDK+CD(MP)*D(L)
0067 CMDK=CMDK+CM(L)*CD(MP)
0068 CMAL=F(KD)*CMAK+CMAL
0069 CML=F(KD)*CMDK+CML
0070 CHA=F(KD)*CHAK+CHA
0071 CHD=F(KD)*CHDK+CHD
0072 XAL=C*AL/(CHA*CT)
0073 XPL=C*ML/(CHD*CT)
0074 XFA=1.0-(PRC-XAL)
0075 XLFD=1.0-(PRC-XDL)
0076 CMAZ=XAL*CLO
0077 CWDZ=XDL*CLO
0078 WRITE(KD,150)Z(1P),XAL,CMAZ,XDL,CWDZ,XLFA,XLFD
0079 1550 FORMAT(20X,F5.2,2(5X,F7.3,4X,FR.4),2(5X,F7.3))
0080 CMINA=CMINA+CMAZ*CT+SIMPSN(1R)
0081 CMIND=CMIND+CWDZ*CT+SIMPSN(1R)
0082 CONTINUE
0083 SA=CMINA/ICLAG*(TW=10.0)
0084 SD=CMIND/ICLGD*(TW=10.0)
0085 CMINA=CMINA/ICTM*10.0
0086 CMIND=CMIND/ICTM*10.0
0087 XLEA=1.0-(XTW-SA)
0088 XLEDM=1.0-(XTW-SD)
0089 WRITE(KG,155)SA,CMIN,SD,CMIND,XLEA,XLEDM
0090 1551 FORMAT(1X,'INTEGRATED',17X,F6.3,5X,F7.4,6X,F6.3,5X,F7.4,2(6X,F6.
0091 13),/)
0092 WRITE(KG,160Q) ZCPA,ZCPD
0093 1600 FORMAT(1X,'SPANWISE POSITIONS OF THE CENTERS OF PRESSURE ARE : ZAP
0094 1,FR.2,5X,17H',FR.2)
0095 WRITE(KG,170S)
0096 1705 FORMAT(1X,'ALPHA',9X,'DELTA',6X,'DEL/ALPH',9X,'CL',12X,'CD',12X
0097 1,'L/1',9X,'CM',10X,'XWL/C',/)
0098 DO 1650 KAL=1,21,5
0099 KAL=KAL-1
0097
```



```

0099      KP=KRA-1
0100      IF(KALA)1500,1500,1505
0101      1500 IF(KR)1501,1501,1507
0102      1501 ALA=0.0
0103      DEL=0.0
0104      DIV=99999.9
0105      GO TO 1510
0106      1502 ALA=0.0
0107      DEL=ELCAT(KP)
0108      DIV=99999.9
0109      GO TO 1510
0110      1505 ALA=ELCAT(KALA)
0111      R=FLNAT(KR)
0112      DEL=ALA+R/10.0
0113      DIV=DEL/ALA
0114      1510 ALPH=ALA/57.296
0115      DELY=DEL/57.296
0116      AA=AA1+ALPH+AA2*DELT
0117      CLALF=CLAG*ALPH
0118      CLDEL=CLNG*DELT
0119      CL=CLALF+CLDEL
0120      SM=J.0
0121      IF(AA)1512,1516,1512
0122      1512 DO 1514 K0=2,KDM
0123      K1=K0*KCM
0124      K2=K0*INDX
0125      1514 SM=SM+12.0*PKN-1.0*((CA(K0)+CA(K1))*ALPH+(CD(K0)+CD(K1)+CD(K2))*
0126      DELT)/AA)*7
0127      CR=(CL+2)*((1.0+SM)/(3.14159*AR)+0.0045*CCNT*CL+7
0128      GO TO 1517
0129      1516 CR=C.0045*CCNT*CL+7
0130      CLD=CL/CR
0131      IF(CL.FD.0.0)XPC=99999.9
0132      IF(CL.FD.0.0)GO TO 1709
0133      CP=(C4THA*CLALF+ALPH+CMIAN*CLDEL+DELT)/CL
0134      XPC=(SA*CLALF+SM*CLDEL)/CL
0135      1709 WRITE(6,1710)MA,DEL,FIV,CL,CR,CLD,CH,XPC
0136      1710 FORMAT(9X,F6.1,7X,F6.1,3X,F6.1,7X,F6.1,5X,F6.1,5X,
0137      F6.4)
0138      1650 CONTINUE
0139      RETURN
0140      END

```

## APPENDIX 4

### Positioning of Control Points Within Their Own Elements

Control points should be located at a point within the element where the effects of the velocities induced by all four boundaries of the element cancel. This is essential in order for the discrete system to converge to the Cauchy-Principal Value of the corresponding continuous singular integral. For a rectangular element this will be, of course, at the centroid of the element.

With the present vortex scheme the elements may be both swept and tapered, so that a question arises as to whether the placement of the control point is critical. A detailed calculation was therefore made for an extreme element with geometric characteristics as shown in Fig. 7.

The velocity induced by the four line vortex elements was computed at finely spaced intervals throughout the interior of the element. In this case, the null point, 0, was found to be displaced from the mid-chord/mid-span point, M, by the amount shown in Fig. 7. This deviation, as well as the velocity at M is extremely small.

In order to obtain an estimate of the error introduced by locating control points at M rather than at 0 in the complete lifting surface program a test run was made with all control points displaced by two percent of their panel chords. This resulted in a change of 0.92% in predicted lift slope. Since this displacement was far in excess of the value shown in Fig. 7, it can be concluded that the error introduced by locating control points at the mid-chord /mid-span position of an element is negligible.

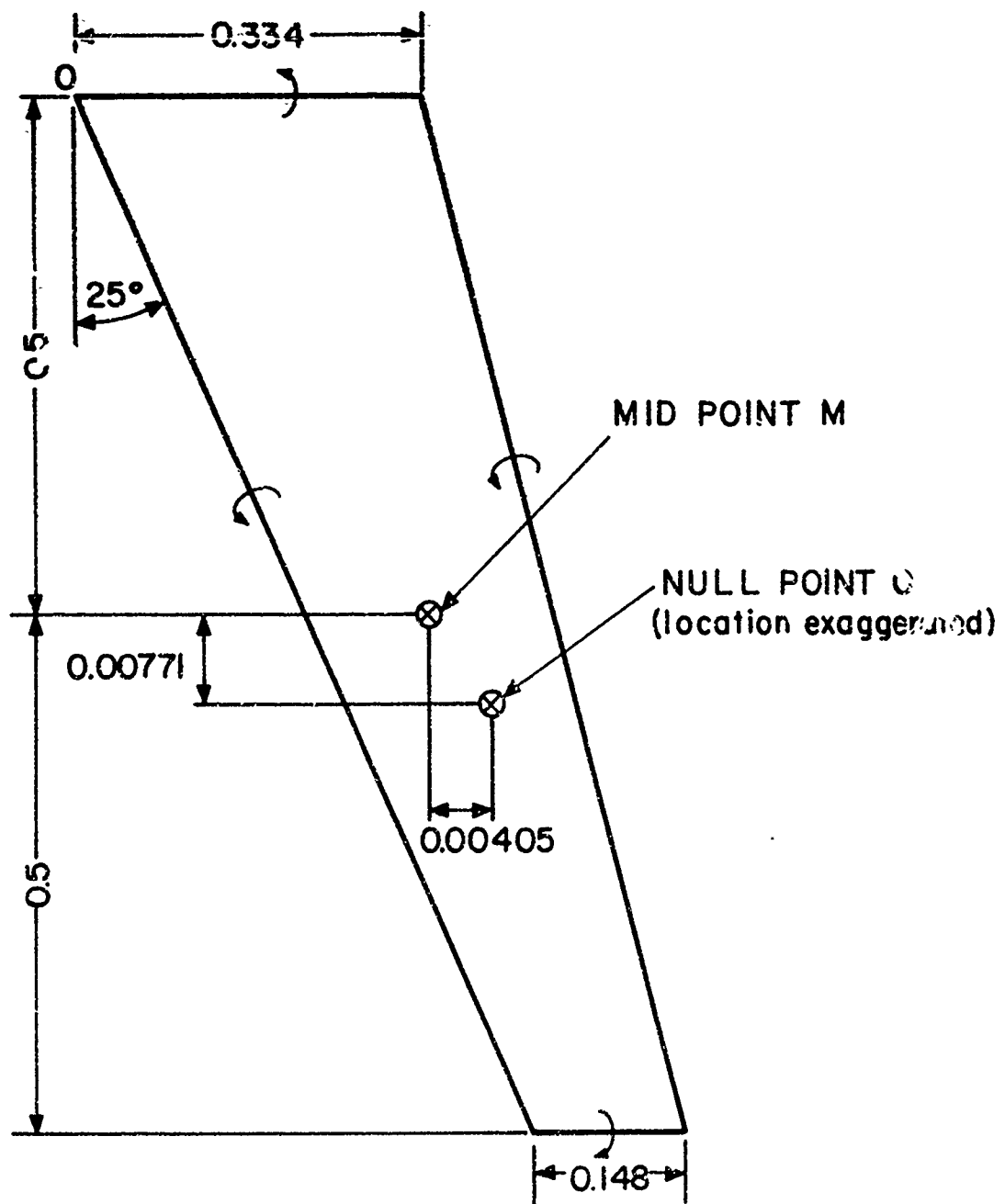


Fig. 7 The Element Used for Estimation of the Error Caused by Control Point Position Deviations